

Linear and nonlinear eigenvalue problems for Dirac systems in unbounded domains^{*†}

Anna Capietto[‡], Walter Dambrosio[‡] and Duccio Papini^{**}

[‡] Dipartimento di Matematica - Università di Torino

Via Carlo Alberto 10, 10123 Torino - Italia

E-mail addresses: anna.capietto@unito.it, walter.dambrosio@unito.it

^{**} Dipartimento di Ingegneria dell'Informazione e Scienze Matematiche - Università di Siena

Via Roma 56, 53100 Siena - Italia

E-mail address: duccio.papini@unisi.it

Abstract

We first study the linear eigenvalue problem for a planar Dirac system in the open half-line and describe the nodal properties of its solution by means of the rotation number. We then give a global bifurcation result for a planar nonlinear Dirac system in the open half-line. As an application, we provide a global continuum of solutions of the nonlinear Dirac equation which have a special form.

Keywords. Dirac system, eigenvalue problem, rotation number, global bifurcation.

1 Introduction

In this paper we give a global bifurcation result (Theorem 4.3) for a nonlinear Dirac system in \mathbb{R}^2 of the form

$$Jz' + P(x)z = \lambda z + S(x, z)z, \quad x > 0, \quad \lambda \in \mathbb{R}, \quad z = (u, v) \in \mathbb{R}^2, \quad (1.1)$$

where

$$J = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

^{*}AMS Subject Classification: 34C23, 34B09, 34L40.

[†]Under the auspices of GNAMPA-I.N.d.A.M., Italy. The work of the first two authors has been performed in the frame of the M.I.U.R. Project 'Topological and Variational Methods in the Study of Nonlinear Phenomena'; the work of the third author has been performed in the frame of the M.I.U.R. Project 'Nonlinear Control: Geometrical Methods and Applications' and of the GNAMPA-I.N.d.A.M. project "Equazione di evoluzione degeneri e singolari: controllo e applicazioni".

and $P(x), S(x, z)$ are continuous symmetric matrices, for every $x > 0$ and $z \in \mathbb{R}^2$. We will be interested in solutions z of (1.1) belonging to the space

$$D_0 = \{z \in L^2(0, +\infty) : z \in AC(0, +\infty), Jz' + P(\cdot)z \in L^2(0, +\infty)\}.$$

In particular, the solutions are convergent to zero at zero and at infinity. This choice is strictly related to the spectral properties of the linear operator $\tau z = Jz' + P(x)z$ and to the possibility of considering self-adjoint extensions of τ (see Section 3).

When P has the form

$$P(x) = P_{V,k,\mu_a}(x) = \begin{pmatrix} -1 + V(x) & -\frac{k}{x} - \mu_a V'(x) \\ -\frac{k}{x} - \mu_a V'(x) & 1 + V(x) \end{pmatrix}, \quad x > 0, \quad (1.2)$$

the differential operator $z \mapsto Jz' + P(\cdot)z$ coincides with the radially symmetric Dirac operator with or without anomalous magnetic moment (cf. [15, 18, 23, 24] and Section 4.2). In this context $V \in C^1(0, +\infty)$ represents an electrostatic potential, $\mu_a \in \mathbb{R}$ an anomalous magnetic moment and $k \in \mathbb{Z} \setminus \{0\}$ (see [23]). For a comprehensive treatment of linear and nonlinear Dirac systems, we refer to the paper by M. Esteban [14]. As it is explained in detail in Section 4.2, nonlinear systems of the form (1.1) arise, for some S , when one is interested in solutions of a nonlinear Dirac PDE which have a special form (cf. (4.49)).

The study of global bifurcation problems for second order equations in unbounded intervals was initiated in the 70s by C. Stuart [22] and N. Dancer [7, 8]. More recent results have been given by P. Rabier-C. Stuart [16], S. Secchi-C. Stuart [20], the first and second author [4] and the authors [5].

In [4] it is considered the particular case when the r.h.s. of (1.1) (and the function S) is regular at zero. We are now able to avoid this restriction and, as a consequence, to treat the physically relevant Dirac operator.

Having in mind a bifurcation result, a comprehensive knowledge of the linear eigenvalue problem

$$Jz' + P(x)z = \lambda z, \quad x > 0, \quad \lambda \in \mathbb{R}, \quad z = (u, v) \in \mathbb{R}^2 \quad (1.3)$$

is necessary. More precisely, we have to study the existence of eigenvalues and their "nodal properties". To this end, in Subsection 2.1, assuming $(\mathcal{P}_1), (\mathcal{P}_2), (\mathcal{P}_3)$ for the matrix P , we first describe (Lemmas 2.6 and 2.13) the behaviour of the solutions of the linear system (1.3) when $x \rightarrow +\infty$ and $x \rightarrow 0^+$. As in [4], we apply the Levinson theorem [13] on the asymptotic properties of solutions of linear equations and, by means of a suitable change of variables, we manage to treat the singularity at zero as well. Using the results of Subsection 2.1, we develop in Subsection 2.2 an oscillatory theory for nontrivial solutions of (2.1) based on the study of the asymptotic behaviour of the angular coordinate θ in the phase-plane (cf. the book by J. Weidmann [24]). It is interesting to observe that, contrary to the case of second order equations, in case of planar Dirac-type systems the angular coordinate is not, in general, an increasing function of x . However, we are able to guarantee (Propositions 2.16 and 2.17) that the limits

$$\theta(+\infty, \lambda) = \lim_{x \rightarrow +\infty} \theta(x, \lambda), \quad \theta(0) = \lim_{x \rightarrow 0^+} \theta(x, \lambda) \quad (1.4)$$

exist and are finite. We can thus give the definition of

$$\text{rot}(z) = \frac{\theta(+\infty, \lambda) - \theta(0)}{\pi}, \quad (1.5)$$

the rotation number of a solution z to (1.3). Roughly speaking, the unboundedness of the interval and the singularity at zero do not prevent solutions to perform only a finite number of rotations around the origin (as in the regular case). A nontrivial phase-plane analysis leads then to some useful continuity properties of the angular function near zero and infinity (Propositions 2.21 and 2.22). In Section 3 we study the spectral theory for the linear operator formally defined by

$$\tau z = Jz' + P(x)z, \quad x > 0. \quad (1.6)$$

More precisely, standard arguments from [24] ensure that τ is in the limit point case at infinity and at zero and that there exists a unique self-adjoint realization A_0 (cf. (3.2)) of τ having (when P has the form (1.2)) essential spectrum $\sigma_{\text{ess}}(A_0) = (-\infty, -1] \cup [1, +\infty)$. Then, the (nontrivial) question of characterizing eigenvalues of A_0 is tackled by the results of Subsection 2.2. Finally, we give results on the existence and accumulation of eigenvalues of A_0 at the boundary of the interval $(-1, 1)$ which are based on the oscillatory behaviour of the solutions for a value of λ corresponding to one of the extrema of the essential spectrum; similar results can be found in the case of second-order differential operators in the book by N. Dunford-J. Schwartz [12] and in case of Dirac operators (without any knowledge of the nodal properties of the corresponding eigenvalues) in the paper by H. Schmid-C. Tretter [18].

Taking advantage of all the results described above, in Subsection 4.1 we give a global bifurcation result (Theorem 4.3) for system (1.1). Due to the fact that we are dealing with an unbounded interval, we face a lack of compactness; this difficulty is overcome by applying an abstract bifurcation result due to C. Stuart [22]. A more precise description of the continuum emanating from eigenvalues of odd multiplicity of the linear operator τ is then performed (as we did in [4]) in Theorem 4.8; to this aim, we develop a continuity-connectivity argument based on a linearization approach and on the properties of the rotation number of a solution to (1.1) (cf. (4.5), (4.33) and Proposition 4.7).

Finally, in Subsection 4.2 we consider the partial differential equation

$$i \sum_{j=1}^3 \alpha_j \frac{\partial \psi}{\partial x_j} - \beta \psi - V(|x|)\psi + ia \sum_{j=1}^3 \alpha_j \frac{\partial V(|x|)}{\partial x_j} \psi = \lambda \psi + \gamma(|x|)F(\langle \beta \psi, \psi \rangle) \beta \psi, \quad x \in \mathbb{R}^3, \quad a \in \mathbb{R}, \quad (1.7)$$

where $\psi : \mathbb{R}^3 \rightarrow \mathbb{C}^4$, $V \in C((0, +\infty), \mathbb{R})$ and $\gamma \in C((0, +\infty), \mathbb{R})$ satisfy suitable assumptions, $\langle \cdot, \cdot \rangle$ denotes the scalar product in \mathbb{C}^4 and α_j ($j = 1, 2, 3$) and β are the 4×4 Dirac matrices (see Subsection 4.2). Set

$$H_0 \psi = i \sum_{j=1}^3 \alpha_j \frac{\partial \psi}{\partial x_j} - \beta \psi, \quad \forall \psi \in H_0^1(\mathbb{R}^3) \subset L^2(\mathbb{R}^3). \quad (1.8)$$

It is well-known (cf. the book by B. Thaller [23]), that there exist suitable subspaces of $L^2(S^2)$ s.t. the restriction of the linear operator $H_0 - V + ia \alpha \cdot \nabla V$ to each of these subspaces can be represented by an ordinary differential operator of the form τ . A remark on the physical meaning of the partial wave subspaces can be found in Remark 4.10 in Section 4.2. It is interesting to observe (on the lines of a paper by F. Cacciafesta [3]) that there are *nonlinear* terms $F(\langle \beta \psi, \psi \rangle) \beta \psi$ in (1.7) which leave the above described subspaces invariant. These appear, among others, in the so-called Soler model and are the most interesting from a physical point of view (cf. [17], [21]). On the same lines, we refer also to the contributions by M. Balabane-T. Cazenave-L. Vazquez [2], Y. Ding-B. Ruf [10], J. Ding, J. Xu, F. Zhang [9], Y. Dong-J. Xie [11] and references therein.

Our contribution (Theorem 4.11) provides the existence of a global continuum of solutions of the nonlinear PDE (1.7) which have a special form (i.e. which belong to one of the above mentioned subspaces). To the authors' knowledge, Theorem 4.11 is the first global bifurcation result for a nonlinear Dirac-type

equation of the form (1.7). In the particular case $V \equiv 0$, M. Balabane-T. Cazenave-A. Douady-F. Merle [1] gave a multiplicity result for solutions (having prescribed nodal properties) to a system of ODEs of the form (4.51). For multiplicity results via critical point theory for the nonlinear Dirac PDE, we refer to Theorem 3.3 in [14] (in case $V \equiv 0$) and to the paper by Y. Ding-B. Ruf [10] (for a potential that includes the Coulomb case). On the other hand, in the particular case of linear Dirac-type systems of ODEs, H. Schmid-C. Tretter [18] have given results for the eigenvalue problem for some special choice of the potential V .

In what follows, we will denote by M_S^2 the set of symmetric 2×2 matrices.

2 Linear Dirac systems

In this Section we consider a linear system of the form

$$Jz' + P(x)z = \lambda z, \quad x > 0, \quad \lambda \in \mathbb{R}, \quad z = (u, v) \in \mathbb{R}^2; \quad (2.1)$$

by a solution of (2.1) we mean a function $z \in AC_{\text{loc}}(0, +\infty)$ satisfying (2.1) almost everywhere in $(0, +\infty)$. In the next Sections we will be interested in solutions $z \in L^2(0, +\infty)$ or $z \in H^1(0, +\infty)$; hence, in describing the solutions of (2.1) we will point out, when possible, if they belong to $L^2(0, +\infty)$ or to $H^1(0, +\infty)$.

We assume that $P \in C((0, +\infty), M_S^2)$ and we denote by p_{ij} its coefficients, as usual. For each pair of real numbers $\mu_- < \mu_+$, let us consider the class \mathcal{P}_μ of continuous maps $P : (0, +\infty) \rightarrow M_S^{2,2}$ satisfying the following conditions:

(\mathcal{P}_1) There exists $q_\infty \geq 1$ such that

$$\lim_{x \rightarrow +\infty} P(x) = \begin{pmatrix} \mu^- & 0 \\ 0 & \mu^+ \end{pmatrix} =: P_\infty \quad (2.2)$$

and

$$\int_1^{+\infty} \|R_\infty(x)\|^{q_\infty} dx < +\infty, \quad (2.3)$$

where $R_\infty(x) = P(x) - P_\infty$, for every $x \geq 1$.

(\mathcal{P}_2) There exist $\beta \geq 1$, $P^* \in M_S^2$ and $q_0 \geq 1$ such that

$$\lim_{x \rightarrow 0^+} x^\beta P(x) = P^* \quad (2.4)$$

and

$$\int_0^1 \frac{1}{x^\beta} \|R_0(x)\|^{q_0} dx < +\infty, \quad (2.5)$$

where $R_0(x) = x^\beta P(x) - P^*$, for every $x \in (0, 1)$.

(\mathcal{P}_3) The matrix P^* satisfies

$$\det P^* < -1/4 \quad \text{if } \beta = 1 \quad (2.6)$$

$$\det P^* < 0 \quad \text{if } \beta > 1. \quad (2.7)$$

In what follows, we write $\Lambda = (\mu^-, \mu^+)$.

Remark 2.1 1. We observe that assumption (\mathcal{P}_2) implies that (2.1) has a singularity for $x \rightarrow 0^+$; indeed, from (2.4) and the fact that P^* is not the zero-matrix (since its determinant is negative in any case), we deduce that

$$p_{ij}(x) \sim \frac{p_{ij}^*}{x^\beta}, \quad x \rightarrow 0^+ \quad (i, j = 1, 2)$$

and, in particular, that $p_{ij} \notin L^1(0, 1)$.

(2) Let us also observe that, for a particular choice of P , the differential operator given in (2.1) coincides with the radially symmetric Dirac operator with or without anomalous magnetic moment (cf. [15, 18, 23, 24] and Section 4.2); indeed, this is the situation when P has the form

$$P(x) = P_{V,k,\mu_a}(x) = \begin{pmatrix} -1 + V(x) & -\frac{k}{x} - \mu_a V'(x) \\ -\frac{k}{x} - \mu_a V'(x) & 1 + V(x) \end{pmatrix}, \quad x > 0, \quad (2.8)$$

where $V \in C^1(0, +\infty)$ is an electrostatic potential, $\mu_a \in \mathbb{R}$ is an anomalous magnetic moment and $k \in \mathbb{Z} \setminus \{0\}$ (see [23]).

Let us assume that V satisfies the following conditions:

$$V(x) = \frac{\gamma_\infty}{x^{\alpha_\infty}} + R_{V,\infty}(x), \quad \alpha_\infty > 0, \quad (2.9)$$

$$x^{\alpha_\infty} R_{V,\infty} = o(1), \quad x^{\alpha_\infty+1} R'_{V,\infty} = o(1), \quad x \rightarrow +\infty$$

and

$$V(x) = \frac{\gamma_0}{x^{\alpha_0}} + R_{V,0}(x), \quad \alpha_0 > 0, \quad (2.10)$$

where

$$\text{if } \mu_a = 0 : \quad \begin{cases} \alpha_0 = 1, & x R_{V,0} = o(1), \quad x \rightarrow 0^+, \\ \int_0^1 \frac{1}{x} |x R_{V,0}(x)|^{q'} dx < +\infty, \quad q' \geq 1, \\ \gamma_0^2 < k^2 - 1/4 \end{cases} \quad (2.11)$$

and

$$\text{if } \mu_a \neq 0 : \quad \begin{cases} x^{\alpha_0} R_{V,0} = o(1), & x^{\alpha_0+1} R'_{V,0} = o(1), \quad x \rightarrow 0^+, \\ \int_0^1 \frac{1}{x^{\alpha_0+1}} |x^{\alpha_0+1} R'_{V,0}(x)|^{q''} dx < +\infty, \quad q'' \geq 1, \\ \gamma_0 \neq 0. \end{cases} \quad (2.12)$$

Under these conditions, assumption (\mathcal{P}_1) is satisfied with $\mu^\pm = \pm 1$; indeed, we obviously have

$$\lim_{x \rightarrow +\infty} P_{V,k,\mu_a}(x) = \lim_{x \rightarrow +\infty} \begin{pmatrix} -1 + V(x) & -\frac{k}{x} - \mu_a V'(x) \\ -\frac{k}{x} - \mu_a V'(x) & 1 + V(x) \end{pmatrix} = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}.$$

Moreover, the matrix R_∞ in (2.3) is given by

$$R_\infty(x) = \begin{pmatrix} V(x) & -\frac{k}{x} - \mu_a V'(x) \\ -\frac{k}{x} - \mu_a V'(x) & V(x) \end{pmatrix}, \quad \forall x > 0;$$

from (2.9) we deduce that $V, R_{V,\infty} \in L^q(1, +\infty)$, for every $q > 1/\alpha_\infty$, and $R'_{V,\infty} \in L^p(1, +\infty)$, for every $p > 1/(\alpha_\infty + 1)$, while we plainly have $k/x \in L^s(1, +\infty)$, for every $s > 1$. By observing that all the functions $V, R_{V,\infty}, R'_{V,\infty}, k/x$ go to zero at infinity, we conclude that (2.3) is satisfied with $q_\infty = 1/(\alpha_\infty + 1)$.

As far as (\mathcal{P}_2) and (\mathcal{P}_3) are concerned, a crucial role is played by the constants γ_0 and μ_a , as it is evident from the assumptions on V . Indeed, let us first discuss the case $\mu_a = 0$; in this situation, taking $\alpha_0 = 1$ and $\beta = \alpha_0 = 1$ in (2.4), we have

$$\begin{aligned} P^* &= \lim_{x \rightarrow 0^+} x P_{V,k,0}(x) = \\ &= \lim_{x \rightarrow 0^+} \begin{pmatrix} -x + xV(x) & -k \\ -k & x + xV(x) \end{pmatrix} = \begin{pmatrix} \gamma_0 & -k \\ -k & \gamma_0 \end{pmatrix} \end{aligned}$$

and

$$R_0(x) = xP(x) - P^* = \begin{pmatrix} -x + xV(x) - \gamma_0 & 0 \\ 0 & x + xV(x) - \gamma_0 \end{pmatrix}.$$

Hence, from (2.10) and (2.11) we infer

$$\int_0^1 \frac{1}{x} \|R_0(x)\|^{q'} dx \leq 2^{q'-1} \int_0^1 (x^{q'-1} + x^{q'-1} |R_{V,0}(x)|^{q'}) dx < +\infty,$$

concluding that (2.5) holds true with $q_0 = q'$. Moreover, the last relation in (2.11) guarantees that (2.6) is fulfilled.

Suppose now $\mu_a \neq 0$; taking $\beta = \alpha_0 + 1$ in (2.4), we have

$$\begin{aligned} P^* &= \lim_{x \rightarrow 0^+} x^{\alpha_0+1} P_{V,k,\mu_a}(x) = \\ &= \lim_{x \rightarrow 0^+} \begin{pmatrix} -x^{\alpha_0+1} + x^{\alpha_0+1} V(x) & -kx^{\alpha_0} - \mu_a x^{\alpha_0+1} V'(x) \\ -kx^{\alpha_0} - \mu_a x^{\alpha_0+1} V'(x) & x^{\alpha_0+1} + x^{\alpha_0+1} V(x) \end{pmatrix} = \begin{pmatrix} 0 & \mu_a \alpha_0 \gamma_0 \\ \mu_a \alpha_0 \gamma_0 & 0 \end{pmatrix} \end{aligned}$$

and

$$R_0(x) = x^{\alpha_0+1}P(x) - P^* = \begin{pmatrix} -x^{\alpha_0+1} + x^{\alpha_0+1}V(x) & -kx^{\alpha_0} - \mu_a x^{\alpha_0+1}V'(x) - \mu_a \alpha_0 \gamma_0 \\ -kx^{\alpha_0} - \mu_a x^{\alpha_0+1}V'(x) - \mu_a \alpha_0 \gamma_0 & x^{\alpha_0+1} + x^{\alpha_0+1}V(x) \end{pmatrix}.$$

Now, let $q_0 > \max(q'', \alpha_0)$; from (2.10) and (2.12) we infer

$$\begin{aligned} \int_0^1 \frac{1}{x^{\alpha_0+1}} |x^{\alpha_0+1}|^{q_0} dx &= \int_0^1 \frac{1}{x^{(\alpha_0+1)(1-q_0)}} dx < +\infty \\ \int_0^1 \frac{1}{x^{\alpha_0+1}} |x^{\alpha_0+1}V(x)|^{q_0} dx &\leq 2^{q_0-1} \int_0^1 \frac{x^{q_0}}{x^{\alpha_0+1}} (\gamma_0^{q_0} + |x^{\alpha_0}R_{V,0}(x)|^{q_0}) dx < +\infty \\ \int_0^1 \frac{1}{x^{\alpha_0+1}} |kx^{\alpha_0}|^{q_0} dx &= |k|^{q_0} \int_0^1 \frac{1}{x^{1+(1-q_0)\alpha_0}} dx < +\infty \\ \int_0^1 \frac{1}{x^{\alpha_0+1}} |-\mu_a x^{\alpha_0+1}V'(x) - \mu_a \alpha_0 \gamma_0|^{q_0} dx &= \\ &= |\mu_a|^{q_0} \int_0^1 \frac{1}{x^{\alpha_0+1}} |x^{\alpha_0+1}R'_{V,0}(x)|^{q''} |x^{\alpha_0+1}R'_{V,0}(x)|^{q_0-q''} dx < +\infty, \end{aligned}$$

concluding again that (2.5) holds true. Moreover, the last relation in (2.12) guarantees that (2.7) is fulfilled.

The fact that $P \in \mathcal{P}_\mu$, and in particular the fact that P satisfies (\mathcal{P}_3) , is related to the spectral properties of the operator $\tau : z \rightarrow Jz' + P(x)z$; indeed, as we will see at the beginning of 3, condition (\mathcal{P}_3) implies that the operator τ is in the limit point case at $x = 0$. As a consequence, it admits a unique self-adjoint realization (cfr. (3.2)); in the particular case of the operator associated to (2.8), with the Coulomb potential

$$V(x) = \frac{\gamma}{x}, \quad \forall x > 0, \quad \gamma < 0, \quad (2.13)$$

condition (\mathcal{P}_3) is satisfied for a larger range of values of γ when $\mu_a \neq 0$. This means that the presence of an anomalous magnetic moment has a regularizing effect on the radial Dirac operator (cf. also [23, Sect. 5.3.2]).

2.1 Asymptotic estimates

In this subsection we describe the behaviour of the solutions of (2.1) when $x \rightarrow +\infty$ or $x \rightarrow 0^+$; this will be the consequence of some general results on the asymptotic properties of solutions of linear equations (see e.g. [13]). As a first step, let us consider a system of the form

$$u' = C(\lambda)u + U(x, \lambda)u, \quad x \geq 1, \lambda \in \Lambda \quad (2.14)$$

where $C(\lambda)$ and $U(x, \lambda)$ are 2×2 matrix, for every $\lambda \in \Lambda$ and $x \geq 1$. We have the following result:

Proposition 2.2 ([13, Th. 1.5.2, Th. 1.8.1, Th. 1.8.2]) *Let us suppose that for every $\lambda \in \Lambda$ the matrix $C(\lambda)$ has two real eigenvalues $\sigma_\lambda^- < 0 < \sigma_\lambda^+$ and let u_λ^- , u_λ^+ be the eigenvectors associated to σ_λ^- and σ_λ^+ , respectively. Moreover, let us assume that*

$$\lim_{x \rightarrow +\infty} U(x, \lambda) = 0, \quad \forall \lambda \in \Lambda, \quad (2.15)$$

and that there exists $q \geq 1$ such that

$$\int_1^{+\infty} \|U(x, \lambda)\|^q dx < +\infty, \quad \forall \lambda \in \Lambda. \quad (2.16)$$

Then, for every $\lambda \in \Lambda$ system (2.14) has two linearly independent solutions $u_{1,\lambda}$ and $u_{2,\lambda}$ satisfying

$$\begin{aligned} u_{1,\lambda}(x) &= (u_\lambda^- + o(1))e^{\sigma_\lambda^-(x-1) + \int_1^x g_{1,\lambda}(t) dt}, \quad x \rightarrow +\infty \\ u_{2,\lambda}(x) &= (u_\lambda^+ + o(1))e^{\sigma_\lambda^+(x-1) + \int_1^x g_{2,\lambda}(t) dt}, \quad x \rightarrow +\infty, \end{aligned} \quad (2.17)$$

where, for $i = 1, 2$, we have

$$\begin{aligned} g_{i,\lambda} &= 0 & \text{if } q = 1 \\ g_{i,\lambda} &\in L^q(1, +\infty) & \text{if } q > 1. \end{aligned} \quad (2.18)$$

PROOF. Let us note that when $q = 1$ the result follows from [13, Th. 1.8.1]. Therefore, assume that $q > 1$; from [13, Th. 1.5.2, Th. 1.8.2] we immediately deduce that (2.17) is satisfied with some functions $g_{i,\lambda}$, $i = 1, 2$, $\lambda \in \Lambda$, such that

$$g_{i,\lambda} = \begin{cases} 0 & \text{if } q = 1 \\ \sum_{m=1}^M g_{i,m,\lambda} & \text{if } q > 1, \end{cases} \quad (2.19)$$

with M such that $2^{M-1} < q \leq 2^M$ and

$$g_{i,m,\lambda} \in L^{q/2^{m-1}}(1, +\infty), \quad \forall m = 1, \dots, M. \quad (2.20)$$

Now, assumption (2.15) implies that

$$\lim_{x \rightarrow +\infty} g_{i,m,\lambda}(x) = 0, \quad \forall m = 1, \dots, M, \lambda \in \Lambda, i = 1, 2$$

(see also formula (1.5.27) in [13]). Hence, for every $i = 1, 2$, $\lambda \in \Lambda$ and $m = 1, \dots, M$ we have

$$g_{i,m,\lambda} \in L^{q/2^{m-1}}(1, +\infty) \Rightarrow g_{i,m,\lambda} \in L^q(1, +\infty).$$

This implies that $g_{i,\lambda} \in L^q(1, +\infty)$, for every $i = 1, 2$ and $\lambda \in \Lambda$. ■

Now, let us observe that

$$f \in L^p(1, +\infty), \quad p > 1 \quad \Rightarrow \quad \left| \int_1^x f(t) dt \right| \leq \|f\|_{L^p} (x-1)^{1/p'}, \quad \forall x \geq 1, \quad (2.21)$$

where p' is the conjugate exponent of p ; noting that in this case $1/p' < 1$, from Proposition 2.2 we obtain the following result:

Proposition 2.3 *Under the assumptions of Proposition 2.2, for every $\lambda \in \Lambda$ we have*

$$\lim_{x \rightarrow +\infty} u_{1,\lambda}(x) = 0 \quad (2.22)$$

and

$$\lim_{x \rightarrow +\infty} \|u_{2,\lambda}(x)\| = +\infty. \quad (2.23)$$

Moreover, if $q > 1$ in (2.16), then there exists $x_1 > 1$ such that

$$\sigma_\lambda^-(x-1) + \int_1^x g_{1,\lambda}(t) dt \leq \frac{\sigma_\lambda^-}{2}(x-1), \quad \forall x \geq x_1. \quad (2.24)$$

Using Proposition 2.2 and Proposition 2.3 we are able to prove some asymptotic results on the solutions of (2.1) when $x \rightarrow +\infty$ or $x \rightarrow 0^+$. We start with the study of (2.1) when $x \rightarrow +\infty$ (cf. also [4]); assume then $x \geq 1$.

Let us first observe that (2.1) can be written as

$$z' = B_\lambda z + Q(x)z, \quad (2.25)$$

where

$$B_\lambda = J^{-1}(\lambda \text{Id} - P_\infty), \quad Q(x) = J^{-1}(P_\infty - P(x)), \quad \forall x > 0.$$

This form corresponds to (2.14) with $C(\lambda) = B_\lambda$ and $U(x, \lambda) = Q(x)$, for every $x \geq 1$, $\lambda \in \Lambda$; note that assumptions (2.2) and (2.3) imply that (2.15) and (2.16), with $q = q_\infty$, hold true. Moreover, if $\lambda \in \Lambda$, setting $\Delta_\lambda = (\mu^+ - \lambda)(\lambda - \mu^-)$, then B_λ has the real eigenvalues $\pm\sqrt{\Delta_\lambda}$; in this situation we denote by $b_{1,\lambda} = (\lambda - \mu^+, \sqrt{\Delta_\lambda})$ and $b_{2,\lambda} = (\mu^+ - \lambda, \sqrt{\Delta_\lambda})$ the eigenvectors of B_λ associated to the eigenvalues $-\sqrt{\Delta_\lambda}$ and $\sqrt{\Delta_\lambda}$, respectively.

Therefore from Proposition 2.2 and Proposition 2.3 we deduce the following results:

Proposition 2.4 *For every $\lambda \in \Lambda$ system (2.1) has two linearly independent solutions $z_{1,\lambda}$ and $z_{2,\lambda}$ satisfying*

$$z_{1,\lambda}(x) = (b_{1,\lambda} + o(1))e^{-\sqrt{\Delta_\lambda}(x-1) + \int_1^x g_1(t) dt}, \quad x \rightarrow +\infty \quad (2.26)$$

$$z_{2,\lambda}(x) = (b_{2,\lambda} + o(1))e^{\sqrt{\Delta_\lambda}(x-1) + \int_1^x g_2(t) dt}, \quad x \rightarrow +\infty,$$

where, for $i = 1, 2$, we have

$$\begin{aligned} g_i &= 0 & \text{if } q_\infty = 1 \\ g_i &\in L^{q_\infty}(1, +\infty) & \text{if } q_\infty > 1. \end{aligned} \quad (2.27)$$

Lemma 2.5 *Assume that $\lambda \in \Lambda$ and let $z_{1,\lambda}$ and $z_{2,\lambda}$ be the solutions of (2.1) given in Proposition 2.4. Then*

$$\lim_{x \rightarrow +\infty} z_{1,\lambda}(x) = 0 \quad (2.28)$$

and

$$\lim_{x \rightarrow +\infty} |(z_{2,\lambda})_1(x)| = \lim_{x \rightarrow +\infty} |(z_{2,\lambda})_2(x)| = +\infty. \quad (2.29)$$

Moreover, $z_{1,\lambda} \in H^1(1, +\infty)$.

PROOF. The relations (2.28) and (2.29) immediately follow from (2.22) and (2.23). In particular (2.29) comes from the fact that neither component of $b_{2,\lambda}$ vanishes.

Moreover, from (2.24) we deduce that there exists $K_{1,\lambda} > 0$ such that

$$\|z_{1,\lambda}(x)\| \leq K_{1,\lambda} e^{-\sqrt{\Delta_\lambda}(x-1)}, \quad \forall x \geq x_1;$$

this implies that $z_{1,\lambda} \in L^2(1, +\infty)$. Now, from the differential equation we deduce that

$$Jz'_{1,\lambda}(x) = \lambda z_{1,\lambda}(x) - P(x)z_{1,\lambda}(x), \quad \forall x \geq 1;$$

since $P \in L^\infty(1, +\infty)$, we infer that $Jz'_{1,\lambda} \in L^2(1, +\infty)$ and then $z_{1,\lambda} \in H^1(1, +\infty)$. ■

Arguing as in the proof of [4, Lemma 2.3], we obtain the following result:

Lemma 2.6 *Assume that $\lambda \in \Lambda$ and let $z = (u, v)$ be a nontrivial solution of (2.1). Then either*

$$\lim_{x \rightarrow +\infty} u(x) = \lim_{x \rightarrow +\infty} v(x) = 0 \tag{2.30}$$

or

$$\lim_{x \rightarrow +\infty} |u(x)| = \lim_{x \rightarrow +\infty} |v(x)| = +\infty. \tag{2.31}$$

Moreover, $z \in H^1(1, +\infty)$ if and only if (2.30) holds true and there exists $\gamma > 0$ such that $z = \gamma z_{1,\lambda}$, where $z_{1,\lambda}$ is given in Proposition 2.4.

Now, let us study the behaviour of the solutions of (2.1) when $x \rightarrow 0^+$; assume then that $x \in (0, 1)$. For every $\beta \geq 1$ let us consider an invertible function $\phi_\beta \in C^1((1, +\infty), (0, 1))$ such that

$$\lim_{t \rightarrow +\infty} \phi_\beta(t) = 0 \quad \text{and} \quad \lim_{t \rightarrow 1^+} \phi_\beta(t) = 1. \tag{2.32}$$

The change of variable $x = \phi_\beta(t)$ transforms (2.1) into

$$w' = -J^{-1}P(\phi_\beta(t))\phi'_\beta(t)w + \lambda J^{-1}\phi'_\beta(t)w, \tag{2.33}$$

where $w(t) = z(\phi_\beta(t))$, for every $t \geq 1$. With a suitable choice of ϕ_β system (2.33) can be reduced to a system of the form (2.14):

Lemma 2.7 *Assume $\beta = 1$ in (\mathcal{P}_2) and let*

$$\phi_\beta(t) = e^{1-t}, \quad \forall t \geq 1.$$

Then (2.33) reduces to a system of the form (2.14) with

$$C = C(\lambda) = J^{-1}P^*, \quad U(t, \lambda) = J^{-1}(R_0(e^{1-t}) - \lambda e^{1-t}\text{Id}), \tag{2.34}$$

for every $t \geq 1$ and $\lambda \in \Lambda$.

Lemma 2.8 Assume $\beta > 1$ in (\mathcal{P}_2) and let

$$\phi_\beta(t) = t^{-1/(\beta-1)}, \quad \forall t \geq 1.$$

Then (2.33) reduces to a system of the form (2.14) with

$$C = C(\lambda) = \frac{1}{\beta-1} J^{-1} P^*, \quad U(t, \lambda) = \frac{1}{\beta-1} J^{-1} (R_0(t^{-1/(\beta-1)}) - \lambda t^{-\beta/(\beta-1)} \text{Id}), \quad (2.35)$$

for every $t \geq 1$ and $\lambda \in \Lambda$.

The proofs of Lemma 2.7 and Lemma 2.8 are straightforward and therefore they are omitted.

Now, set $\Delta^* = -\det P^*$ and observe that the matrix C given in (2.34) or (2.35) has the eigenvalues $\sigma^\pm = \pm\sqrt{\Delta^*}$ if $\beta = 1$ and $\sigma^\pm = \pm\sqrt{\Delta^*}/(\beta-1)$ if $\beta > 1$; in what follows, we will denote by w_1^* and w_2^* the eigenvectors of C associated to σ^\pm .

Moreover, from (2.4) and the definition of R_0 we deduce that the function U given in (2.34) or (2.35) satisfies (2.15). Finally, let us note that (2.5) implies that (2.16) is satisfied with $q = q_0$; indeed, when $\beta = 1$ we have

$$\begin{aligned} \int_1^{+\infty} \|U(t, \lambda)\|^{q_0} dt &\leq \left\{ \left[\int_1^{+\infty} \|R_0(e^{1-t})\|^{q_0} dt \right]^{\frac{1}{q_0}} + \lambda \left[\int_1^{+\infty} e^{q_0(1-t)} dt \right]^{\frac{1}{q_0}} \right\}^{q_0} = \\ &= \left\{ \left[\int_0^1 \frac{1}{x} \|R_0(x)\|^{q_0} dx \right]^{\frac{1}{q_0}} + \frac{\lambda}{q_0^{1/q_0}} \right\}^{q_0} < +\infty. \end{aligned}$$

On the other hand, if $\beta > 1$ we deduce that

$$\begin{aligned} \int_1^{+\infty} \|U(t, \lambda)\|^{q_0} dt &\leq \frac{1}{(\beta-1)^{q_0}} \left\{ \left[\int_1^{+\infty} \|R_0(t^{-1/(\beta-1)})\|^{q_0} dt \right]^{\frac{1}{q_0}} + \lambda \left[\int_1^{+\infty} t^{-\beta q_0/(\beta-1)} dt \right]^{\frac{1}{q_0}} \right\}^{q_0} = \\ &= \frac{1}{\beta-1} \left\{ \left[(\beta-1) \int_0^1 \frac{1}{x^\beta} \|R_0(x)\|^{q_0} dx \right]^{\frac{1}{q_0}} + \lambda \left(\frac{\beta-1}{q_0} \right)^{\frac{1}{q_0}} \right\}^{q_0} < +\infty. \end{aligned}$$

Therefore, we can apply Proposition 2.2 and Proposition 2.3 to (2.33), with ϕ_β as above, and obtain the following results:

Proposition 2.9 For every $\lambda \in \Lambda$ system (2.33), with ϕ_β as in Lemma 2.7 or Lemma 2.8, has two linearly independent solutions $w_{1,\lambda}$ and $w_{2,\lambda}$ satisfying

$$\begin{aligned} w_{1,\lambda}(t) &= (w_1^* + o(1)) e^{-\sqrt{\Delta^*}(t-1) + \int_1^t g_{1,\lambda}(s) ds}, \quad t \rightarrow +\infty \\ w_{2,\lambda}(t) &= (w_2^* + o(1)) e^{\sqrt{\Delta^*}(t-1) + \int_1^t g_{2,\lambda}(s) ds}, \quad t \rightarrow +\infty, \end{aligned} \quad (2.36)$$

if $\beta = 1$ and

$$\begin{aligned} w_{1,\lambda}(t) &= (w_1^* + o(1))e^{-\frac{\sqrt{\Delta^*}}{\beta-1}(t-1) + \int_1^t g_{1,\lambda}(s) ds}, \quad t \rightarrow +\infty \\ w_{2,\lambda}(t) &= (w_2^* + o(1))e^{\frac{\sqrt{\Delta^*}}{\beta-1}(t-1) + \int_1^t g_{2,\lambda}(s) ds}, \quad t \rightarrow +\infty, \end{aligned} \quad (2.37)$$

if $\beta > 1$, where, for $i = 1, 2$, we have

$$\begin{aligned} g_{i,\lambda} &= 0 & \text{if } q_0 = 1 \\ g_{i,\lambda} &\in L^{q_0}(1, +\infty) & \text{if } q_0 > 1. \end{aligned} \quad (2.38)$$

Lemma 2.10 Assume that $\lambda \in \Lambda$ and let $w_{1,\lambda}$ and $w_{2,\lambda}$ be the solutions of (2.33) given in Proposition 2.9. Then

$$\lim_{t \rightarrow +\infty} w_{1,\lambda}(t) = 0 \quad (2.39)$$

and

$$\lim_{t \rightarrow +\infty} \|w_{2,\lambda}(t)\| = +\infty. \quad (2.40)$$

Moreover, the solution $w_{1,\lambda}$ satisfies

$$\int_1^{+\infty} \|w_{1,\lambda}(t)\|^2 e^t dt < +\infty, \quad \text{if } \beta = 1 \quad (2.41)$$

and

$$\int_1^{+\infty} \|w_{1,\lambda}(t)\|^2 t^{\beta/(\beta-1)} dt < +\infty, \quad \text{if } \beta > 1. \quad (2.42)$$

PROOF. Let us note that (2.39) and (2.40) immediately follow from (2.22) and (2.23).

As far as (2.41) is concerned, from (2.36) we deduce that there exists $K_{1,\lambda} > 0$ such that

$$\|w_{1,\lambda}(t)\|^2 e^t \sim K_{1,\lambda} e^{-2\sqrt{\Delta^*}(t-1) + 2\int_1^t g_{1,\lambda}(s) ds} e^t, \quad t \rightarrow +\infty; \quad (2.43)$$

now, let us observe that $1 - 2\sqrt{\Delta^*} > 0$, since (2.6) holds. Hence, using again (2.21) we infer that there exists $t_1 > 1$ such that

$$e^{(1-2\sqrt{\Delta^*})t + 2\int_1^t g_{1,\lambda}(s) ds} \leq e^{(1-2\sqrt{\Delta^*})t/2}, \quad \forall t \geq t_1, \quad (2.44)$$

is satisfied. Conditions (2.43) and (2.44) imply (2.41).

Finally, when $\beta > 1$ from (2.37) we deduce that there exists $M_{1,\lambda} > 0$ such that

$$\|w_{1,\lambda}(t)\|^2 t^{\beta/(\beta-1)} \sim M_{1,\lambda} e^{-2\frac{\sqrt{\Delta^*}}{\beta-1}(t-1) + 2\int_1^t g_{1,\lambda}(s) ds} t^{\beta/(\beta-1)}, \quad t \rightarrow +\infty; \quad (2.45)$$

moreover, from (2.24) we infer that there exists $t_2 > 1$ such that

$$e^{-2\frac{\sqrt{\Delta^*}}{\beta-1}(t-1) + 2\int_1^t g_{1,\lambda}(s) ds} \leq e^{-\frac{\sqrt{\Delta^*}}{\beta-1}(t-1)}, \quad \forall t \geq t_2, \quad (2.46)$$

is satisfied. Conditions (2.45) and (2.46) imply (2.42). ■

The next result is a consequence of Proposition 2.9 and Lemma 2.10.

Proposition 2.11 *For every $\lambda \in \Lambda$ system (2.1) has two linearly independent solutions $\zeta_{1,\lambda}$ and $\zeta_{2,\lambda}$ satisfying*

$$\begin{aligned}\zeta_{1,\lambda}(x) &= (w_1^* + o(1)) x^{\sqrt{\Delta^*}} e^{\int_1^{1-\log x} g_{1,\lambda}(s) ds}, \quad x \rightarrow 0^+ \\ \zeta_{2,\lambda}(x) &= (w_2^* + o(1)) x^{-\sqrt{\Delta^*}} e^{\int_1^{1-\log x} g_{2,\lambda}(s) ds}, \quad x \rightarrow 0^+, \end{aligned} \quad (2.47)$$

if $\beta = 1$ and

$$\begin{aligned}\zeta_{1,\lambda}(x) &= (w_1^* + o(1)) e^{-\frac{\sqrt{\Delta^*}}{\beta-1} x^{1-\beta} + \int_1^{-\log x} g_{1,\lambda}(s) ds}, \quad x \rightarrow 0^+ \\ \zeta_{2,\lambda}(x) &= (w_2^* + o(1)) e^{\frac{\sqrt{\Delta^*}}{\beta-1} x^{1-\beta} + \int_1^{-\log x} g_{2,\lambda}(s) ds}, \quad x \rightarrow 0^+, \end{aligned} \quad (2.48)$$

if $\beta > 1$, where, for $i = 1, 2$, we have

$$\begin{aligned}g_{i,\lambda} &= 0 & \text{if } q_0 = 1 \\ g_{i,\lambda} &\in L^{q_0}(1, +\infty) & \text{if } q_0 > 1. \end{aligned} \quad (2.49)$$

Lemma 2.12 *Assume that $\lambda \in \Lambda$ and let $\zeta_{1,\lambda}$ and $\zeta_{2,\lambda}$ be the solutions of (2.1) given in Proposition 2.11. Then*

$$\lim_{x \rightarrow 0^+} \zeta_{1,\lambda}(x) = 0 \quad (2.50)$$

and

$$\lim_{x \rightarrow 0^+} \|\zeta_{2,\lambda}(x)\| = +\infty. \quad (2.51)$$

Moreover, $\zeta_{1,\lambda} \in H^1(0, 1)$.

PROOF. The relations (2.50) and (2.51) immediately follow from (2.39) and (2.40).

Now, assume that $\beta = 1$; let us observe that we have

$$\int_0^1 \frac{\|\zeta_{1,\lambda}(x)\|^2}{x^2} dx = \int_1^{+\infty} \|\zeta_{1,\lambda}(e^{1-t})\|^2 e^{t-1} dt = \int_1^{+\infty} \|w_{1,\lambda}(t)\|^2 e^{t-1} dt < +\infty, \quad (2.52)$$

by (2.41). This condition obviously implies that

$$\int_0^1 \|\zeta_{1,\lambda}(x)\|^2 dx \leq \int_0^1 \frac{\|\zeta_{1,\lambda}(x)\|^2}{x^2} dx < +\infty \quad (2.53)$$

and so $\zeta_{1,\lambda} \in L^2(0, 1)$; in order to prove that $\zeta'_{1,\lambda} \in L^2(0, 1)$, let us note that $\zeta_{1,\lambda}$ satisfies

$$J\zeta'_{1,\lambda}(x) = \lambda\zeta_{1,\lambda}(x) - P(x)\zeta_{1,\lambda}(x), \quad \forall x \in (0, 1).$$

From (\mathcal{P}_2) we deduce that there exists $K > 0$ such that

$$\|P(x)\| \leq \frac{K}{x}, \quad \forall x \in (0, 1);$$

therefore, (2.52) implies that $P\zeta_{1,\lambda} \in L^2(0, 1)$ and then also $\zeta'_{1,\lambda} \in L^2(0, 1)$.

When $\beta > 1$ by (2.42) we obtain

$$\int_0^1 \frac{\|\zeta_{1,\lambda}(x)\|^2}{x^{2\beta}} dx = \frac{1}{\beta-1} \int_1^{+\infty} \frac{\|\zeta_{1,\lambda}(t^{-1/(\beta-1)})\|^2}{t^{-2\beta/(\beta-1)}} t^{-\beta/(\beta-1)} dt = \int_1^{+\infty} \|w_{1,\lambda}(t)\|^2 t^{\beta/(\beta-1)} dt < +\infty. \quad (2.54)$$

This condition implies that

$$\int_0^1 \|\zeta_{1,\lambda}(x)\|^2 dx \leq \int_0^1 \frac{\|\zeta_{1,\lambda}(x)\|^2}{x^{2\beta}} dx < +\infty \quad (2.55)$$

and so $\zeta_{1,\lambda} \in L^2(0, 1)$; arguing as above, in order to prove that $\zeta'_{1,\lambda} \in L^2(0, 1)$, let us note that $\zeta_{1,\lambda}$ satisfies

$$J\zeta'_{1,\lambda}(x) = \lambda\zeta_{1,\lambda}(x) - P(x)\zeta_{1,\lambda}(x), \quad \forall x \in (0, 1).$$

From (\mathcal{P}_2) we deduce that there exists $M > 0$ such that

$$\|P(x)\| \leq \frac{M}{x^\beta}, \quad \forall x \in (0, 1);$$

therefore, (2.54) implies that $P\zeta_{1,\lambda} \in L^2(0, 1)$ and then also $\zeta'_{1,\lambda} \in L^2(0, 1)$. ■

Lemma 2.13 *Assume that $\lambda \in \Lambda$ and let $z = (u, v)$ be a nontrivial solution of (2.1). Then either*

$$\lim_{x \rightarrow 0^+} u(x) = \lim_{x \rightarrow 0^+} v(x) = 0 \quad (2.56)$$

or

$$\lim_{x \rightarrow 0^+} \|z(x)\| = +\infty. \quad (2.57)$$

Moreover, $z \in H^1(0, 1)$ if and only if (2.56) holds true and there exists $\xi \in \mathbb{R}$ such that $z = \xi\zeta_{1,\lambda}$, where $\zeta_{1,\lambda}$ is given in Proposition 2.11.

Remark 2.14 *Let us denote by Z the set of solutions of (2.1). From Proposition 2.4 and Proposition 2.11 we deduce that*

$$Z = \text{span} \{z_{1,\lambda}, z_{2,\lambda}\} = \text{span} \{\zeta_{1,\lambda}, \zeta_{2,\lambda}\}.$$

As far as nontrivial solutions $z \in L^2(0, +\infty)$ are concerned, let us observe that Lemma 2.6 and Lemma 2.13 prove that

$$z \in L^2(1, +\infty) \iff z \in \text{span} \{z_{1,\lambda}\} := Z_\infty$$

and

$$z \in L^2(0, 1) \iff z \in \text{span} \{\zeta_{1,\lambda}\} := Z_0.$$

As a consequence, $z \in L^2(0, +\infty)$ is a solution of (2.1) if and only if

$$z \in Z_0 \cap Z_\infty.$$

We conclude this subsection with some explicit formulas for solutions of the non-homogeneous equation

$$Jz' + P(x)z = f, \quad (2.58)$$

where $z, f \in L^2(0, 1)$. They are based on the fact that the homogeneous equation (2.1) has a suitable dichotomy at zero when $\lambda = 0$ (see [6]).

First of all, let us observe that from [6, §3] we deduce that (2.58) has a solution $z_f \in L^\infty(0, 1)$ when $f \in L^2(0, 1)$. Moreover, let us point out that the previous results on the asymptotic behaviour for $x \rightarrow 0^+$ of the solutions of (2.1) hold true also when $\lambda = 0$; indeed, they are based on the fact that $\Delta^* > 0$. Hence, according to Remark 2.14, all the solutions $z \in L^2(0, 1)$ of (2.58) are of the form

$$z = c\zeta_{1,0} + z_f,$$

for some $c \in \mathbb{R}$. More precisely, we have the following result:

Theorem 2.15 ([6, §3]) *Let us consider $f \in L^2(0, 1)$ and let $z \in L^2(0, 1)$ be a solution of (2.58). Then, there exist $c \in \mathbb{R}$ and $G : (0, 1) \times (0, 1) \rightarrow \mathbb{R}^2$ such that*

$$z(x) = c\zeta_{1,0}(x) + \int_0^1 G(x, \xi) f(\xi) d\xi, \quad \forall x \in (0, 1). \quad (2.59)$$

Moreover, there exist $K > 0$ such that

$$\|G(x, \xi)\| \leq \begin{cases} K \left(\frac{\min(x, \xi)}{\max(x, \xi)} \right)^{\sqrt{\Delta^*}} & \text{if } \beta = 1 \\ K \left(\frac{e^{-\min(x, \xi)^{1-\beta}}}{e^{-\max(x, \xi)^{1-\beta}}} \right)^{\sqrt{\Delta^*}/(\beta-1)} & \text{if } \beta > 1 \end{cases} \leq K, \quad \forall (x, \xi) \in (0, 1) \times (0, 1). \quad (2.60)$$

PROOF. We just point out that the result follows from the change of variables $x = \phi_\beta(t)$, from estimates in [6, §3, formulas (3) and (4)] and Propositions 2.9 or 2.11. \blacksquare

2.2 Oscillatory properties

In this subsection we develop an oscillatory theory for nontrivial solutions of (2.1), based on the study of the angular coordinate in the phase-plane (see [24]). For every nontrivial solution (u, v, λ) of (2.1) let us introduce the polar coordinates $(\rho, \theta) = (\rho(x, \lambda), \theta(x, \lambda))$ according to

$$\begin{cases} u = \rho \cos \theta \\ v = \rho \sin \theta. \end{cases}$$

Observe that θ is defined mod. 2π ; we do not impose a normalization condition on θ and then the following results hold true for any angular coordinate associated to a nontrivial solution z . As a first step, we are able to study the asymptotic behaviour of θ when $x \rightarrow +\infty$ or $x \rightarrow 0^+$; this follows from the results of Subsection 2.1.

Proposition 2.16 ([4, Prop. 2.4]) *For every $\lambda \in \Lambda$ the function $\theta(\cdot, \lambda)$ has limit at infinity and we have either*

$$\lim_{x \rightarrow +\infty} \theta(x, \lambda) = \pi - \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}} \pmod{\pi} \quad (2.61)$$

or

$$\lim_{x \rightarrow +\infty} \theta(x, \lambda) = \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}} \pmod{\pi}. \quad (2.62)$$

Moreover, (2.61) and (2.62) correspond to the cases when (2.30) and (2.31) are fulfilled, respectively.

Proposition 2.17 *For every $\lambda \in \Lambda$ the function $\theta(\cdot, \lambda)$ has limit at zero and we have either*

$$\lim_{x \rightarrow 0^+} \theta(x, \lambda) = \arctan \frac{w_{1,2}^*}{w_{1,1}^*} \pmod{\pi} \quad (2.63)$$

or

$$\lim_{x \rightarrow 0^+} \theta(x, \lambda) = \arctan \frac{w_{2,2}^*}{w_{2,1}^*} \pmod{\pi}, \quad (2.64)$$

where w_1^* and w_2^* are eigenvectors of C associated to the eigenvalues σ^\pm , respectively. Moreover, (2.63) and (2.64) correspond to the cases when (2.56) and (2.57) are fulfilled, respectively.

Let us observe that the possible limits of $\theta(\cdot, \lambda)$ at zero do not depend on $\lambda \in \Lambda$; in what follows, we denote

$$\theta(+\infty, \lambda) = \lim_{x \rightarrow +\infty} \theta(x, \lambda)$$

$$\theta(0) = \lim_{x \rightarrow 0^+} \theta(x, \lambda),$$

which exist and are finite by Proposition 2.16 and Proposition 2.17.

Remark 2.18 *According to Remark 2.14 and the above Propositions, we deduce that for a nontrivial solution z of (2.1) we have*

$$z \in L^2(1, +\infty) \iff \theta(+\infty, \lambda) = \pi - \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}} \pmod{\pi}$$

and

$$z \in L^2(0, 1) \iff \theta(0) = \arctan \frac{w_{1,2}^*}{w_{1,1}^*} \pmod{\pi}.$$

Proposition 2.16 and Proposition 2.17 imply that any angular function $\theta(\cdot, \lambda)$ is bounded on $(0, +\infty)$, for every $\lambda \in \Lambda$. As a consequence, we can associate to every nontrivial solution z of (2.1) the rotation number

$$\text{rot}(z) = \frac{\theta(+\infty, \lambda) - \theta(0)}{\pi}. \quad (2.65)$$

Roughly speaking, the unboundedness of the interval and the singularity at zero do not prevent solutions to perform only a finite number of rotations around the origin (as in the regular case). It is important to observe that $\text{rot}(z)$ does not depend on the choice of the angular function of z . In Section 4 we will study some continuity properties of the rotation number defined in (2.65).

We conclude this subsection with some asymptotic phase-plane analysis for (2.1); as above, we prove the results for $x \rightarrow +\infty$. The case of $x \rightarrow 0^+$ can be obtained in an analogous way by means of the change of variable $x = \phi_\beta(t)$ already introduced.

Let us consider again (2.25), which is equivalent to (2.1), and a similar system

$$Jz' + \tilde{P}(x)z = \tilde{\lambda}z, \quad (2.66)$$

where $\tilde{P} \in \mathcal{P}_\mu$ and $\tilde{\lambda} \in \Lambda$; (2.66) can be written in the form

$$z' = B_{\tilde{\lambda}}z + \tilde{Q}(x)z, \quad (2.67)$$

where $B_{\tilde{\lambda}} = J^{-1}(\tilde{\lambda}\text{Id} - P_\infty)$ and $\tilde{Q}(x) = J^{-1}(P_\infty - \tilde{P}(x))$, for every $x > 0$. Let us note that the matrix P_∞ is the same both for P and \tilde{P} , since $P, \tilde{P} \in \mathcal{P}_\mu$.

For every $\lambda \in \Lambda$, let $b_{1,\lambda}, b_{2,\lambda}$ be as in Proposition 2.4; from the discussion leading to Proposition 2.4 we know that

$$b_{1,\lambda} = (\lambda - \mu^+, \sqrt{\Delta_\lambda}), \quad b_{2,\lambda} = (-\lambda + \mu^+, \sqrt{\Delta_\lambda}),$$

for every $\lambda \in \Lambda$; moreover, there exists $\rho_\lambda > 0$ such that

$$b_{1,\lambda} = \rho_\lambda(\cos \theta_{\infty,\lambda}, \sin \theta_{\infty,\lambda})$$

$$b_{2,\lambda} = \rho_\lambda(-\cos \theta_{\infty,\lambda}, \sin \theta_{\infty,\lambda}),$$

where

$$\theta_{\infty,\lambda} = \pi - \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}} \pmod{\pi}.$$

For every $\theta \in \mathbb{R}$, let r_θ be the straight line of equation $x \sin \theta - y \cos \theta = 0$ and let $v_\theta = (\sin \theta, -\cos \theta)$; moreover, let r_θ^\pm be the half-lines given by the intersection of r_θ with the half-planes $H^+ = \{(x, y) \in \mathbb{R}^2 : x > 0\}$ and $H^- = \{(x, y) \in \mathbb{R}^2 : x < 0\}$, respectively. We are in position to prove the following result:

Proposition 2.19 *For every $\tilde{\lambda} \in \Lambda$, $\tilde{P} \in \mathcal{P}_\mu$ and for every $\theta \in (\pi/2, \pi)$ there exist $\tilde{\delta} > 0$ and $\tilde{x}_\infty = \tilde{x}_\infty(\tilde{\lambda}, \tilde{P}, \theta) > 0$ such that for every $\lambda \in \Lambda$ and $P \in \mathcal{P}_\mu$ with*

$$|\lambda - \tilde{\lambda}| < \tilde{\delta}, \quad \|P - \tilde{P}\|_{L^\infty(1, +\infty)} < \tilde{\delta} \quad (2.68)$$

we have

$$\begin{aligned} \theta < \theta_{\infty,\lambda} &\Rightarrow \langle v_\theta, B_\lambda w + Q(x)w \rangle > 0, \quad \forall w \in r_\theta^-, \quad \forall x \geq \tilde{x}_\infty \\ \theta > \theta_{\infty,\lambda} &\Rightarrow \langle v_\theta, B_\lambda w + Q(x)w \rangle < 0, \quad \forall w \in r_\theta^-, \quad \forall x \geq \tilde{x}_\infty. \end{aligned} \quad (2.69)$$

PROOF. First of all, let us observe that it is sufficient to prove (2.69) when w is a versor. Therefore, let $w = (\cos \theta, \sin \theta)$; a simple computation shows that

$$\phi_\lambda(\theta, w) := \langle v_\theta, B_\lambda w \rangle = \cos^2 \theta ((\mu^+ - \lambda) \tan^2 \theta - (\lambda - \mu^-)), \quad \forall \lambda \in \Lambda. \quad (2.70)$$

Let us fix $\tilde{\lambda} \in \Lambda$, $\tilde{P} \in \mathcal{P}_\mu$ and $\theta \in (\pi/2, \pi)$ such that $\theta < \theta_{\infty, \tilde{\lambda}}$; the continuity of $\theta_{\infty, \lambda}$ as a function of $\lambda \in \Lambda$ implies that there exists $\delta_1 > 0$ such that $\theta < \theta_{\infty, \lambda}$ if $|\lambda - \tilde{\lambda}| < \delta_1$.

From (2.70) we deduce that

$$\begin{aligned} \phi_{\tilde{\lambda}}(\theta_{\infty, \tilde{\lambda}}, w) &= 0 \\ \theta < \theta_{\infty, \tilde{\lambda}} &\Rightarrow \phi_{\tilde{\lambda}}(\theta, w) > \phi_{\tilde{\lambda}}(\theta_{\infty, \tilde{\lambda}}, w) = 0. \end{aligned}$$

Hence, there exists $\delta_2 \in (0, \delta_1)$ such that

$$|\lambda - \tilde{\lambda}| < \delta_2 \Rightarrow \theta < \theta_{\infty, \lambda} \quad \text{and} \quad \phi_\lambda(\theta, w) > \frac{\phi_{\tilde{\lambda}}(\theta, w)}{4} > 0 \quad (2.71)$$

Now, from assumption (2.2) we deduce that

$$\lim_{x \rightarrow +\infty} \langle v_\theta, \tilde{Q}(x)w \rangle = \lim_{x \rightarrow +\infty} \langle v_\theta, J^{-1}(P_\infty - \tilde{P}(x))w \rangle = 0;$$

this implies that there exists $\tilde{x}_\infty = \tilde{x}_\infty(\tilde{\lambda}, \tilde{P}, \theta) > 1$ such that

$$x \geq \tilde{x}_\infty \Rightarrow |\langle v_\theta, \tilde{Q}(x)w \rangle| < \frac{\phi_{\tilde{\lambda}}(\theta, w)}{16}. \quad (2.72)$$

On the other hand, setting $\delta_3 = \phi_{\tilde{\lambda}}(\theta, w)/16$, if $\|P - \tilde{P}\|_{L^\infty(1, +\infty)} < \delta_3$ we have

$$\begin{aligned} |\langle v_\theta, Q(x)w \rangle - \langle v_\theta, \tilde{Q}(x)w \rangle| &= |\langle v_\theta, J^{-1}(\tilde{P}(x) - P(x))w \rangle| \leq \\ &\leq \|\tilde{P}(x) - P(x)\| < \frac{\phi_{\tilde{\lambda}}(\theta, w)}{16}, \quad \forall x \geq 1. \end{aligned} \quad (2.73)$$

From (2.72) and (2.73) we deduce that

$$\|P - \tilde{P}\|_{L^\infty(1, +\infty)} < \delta_3, \quad x \geq \tilde{x}_\infty \Rightarrow |\langle v_\theta, Q(x)w \rangle| < \frac{\phi_{\tilde{\lambda}}(\theta, w)}{8}. \quad (2.74)$$

Now, let us set $\tilde{\delta} = \min(\delta_2, \delta_3)$; when $|\lambda - \tilde{\lambda}| < \tilde{\delta}$ and $\|P - \tilde{P}\|_{L^\infty(1, +\infty)} < \tilde{\delta}$ both (2.71) and (2.74) hold true. As a consequence, we obtain

$$x \geq \tilde{x}_\infty \Rightarrow \langle v_\theta, B_\lambda w + Q(x)w \rangle > \frac{\phi_{\tilde{\lambda}}(\theta, w)}{4} - \frac{\phi_{\tilde{\lambda}}(\theta, w)}{8} > 0, \quad (2.75)$$

i.e. the first inequality in (2.69) is satisfied.

An analogous argument proves the validity of the second inequality in (2.69). ■

In a very similar way it is possible to prove the following Proposition:

Proposition 2.20 For every $\tilde{\lambda} \in \Lambda$, $\tilde{P} \in \mathcal{P}_\mu$ and for every $\theta \in (0, \pi/2)$ there exist $\tilde{\delta}_1 > 0$ and $\tilde{x}_{\infty,1} = \tilde{x}_{\infty,1}(\tilde{\lambda}, \tilde{P}, \theta) > 0$ such that for every $\lambda \in \Lambda$ and $P \in \mathcal{P}_\mu$ with

$$|\lambda - \tilde{\lambda}| < \tilde{\delta}, \quad \|P - \tilde{P}\|_{L^\infty(1, +\infty)} < \tilde{\delta} \quad (2.76)$$

we have

$$\begin{aligned} \theta < \pi - \theta_{\infty, \lambda} &\Rightarrow \langle v_\theta, B_\lambda w + Q(x)w \rangle < 0, \quad \forall w \in r_\theta^+, \quad \forall x \geq \tilde{x}_{\infty,1} \\ \theta > \pi - \theta_{\infty, \lambda} &\Rightarrow \langle v_\theta, B_\lambda w + Q(x)w \rangle > 0, \quad \forall w \in r_\theta^+, \quad \forall x \geq \tilde{x}_{\infty,1}. \end{aligned} \quad (2.77)$$

From Proposition 2.19 and Proposition 2.20 we deduce the following result:

Proposition 2.21 For every $\tilde{\lambda} \in \Lambda$ and $\tilde{P} \in \mathcal{P}_\mu$ there exists $\tilde{\epsilon} > 0$ such that for every $\epsilon \in (0, \tilde{\epsilon})$ there exist $\tilde{\delta} > 0$ and $\tilde{x}_\infty = \tilde{x}_\infty(\tilde{\lambda}, \tilde{P}, \epsilon) > 0$ such that for every $\lambda \in \Lambda$ and $P \in \mathcal{P}_\mu$ with

$$|\lambda - \tilde{\lambda}| < \tilde{\delta}, \quad \|P - \tilde{P}\|_{L^\infty(1, +\infty)} < \tilde{\delta} \quad (2.78)$$

and for every nontrivial solution $z \in L^2((1, +\infty))$ of (2.1) we have

$$|\theta(x, \lambda) - \theta(+\infty, \lambda)| < \epsilon, \quad \forall x \geq \tilde{x}_\infty, \quad (2.79)$$

where $\theta(\cdot, \lambda)$ is any angular coordinate of z .

PROOF. Without loss of generality let us assume that

$$\theta_{\infty, \lambda} = \pi - \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}} \in \left(\frac{\pi}{2}, \pi\right)$$

and define

$$\tilde{\epsilon} = \min \left\{ \theta_{\infty, \tilde{\lambda}} - \frac{\pi}{2}, \pi - \theta_{\infty, \tilde{\lambda}} \right\} > 0.$$

Fix any $\epsilon \in (0, \tilde{\epsilon})$ and consider

$$\theta_1 = \theta_{\infty, \tilde{\lambda}} - \frac{\epsilon}{2} \quad \text{and} \quad \theta_2 = \theta_{\infty, \tilde{\lambda}} + \frac{\epsilon}{2},$$

thus the cone between r_{θ_1} and r_{θ_2} lies inside the II and the IV quadrants and its angular amplitude is exactly ϵ . We use the continuity of $\theta_{\infty, \lambda}$ with respect to λ and apply Proposition 2.19 twice with the choices $\theta = \theta_1$ and $\theta = \theta_2$ in order to find $\tilde{\delta} > 0$ and $\tilde{x}_\infty = \tilde{x}_\infty(\tilde{\lambda}, \tilde{P}, \epsilon) > 0$ in such a way that, if (2.78) hold, then $|\theta_{\infty, \lambda} - \theta_{\infty, \tilde{\lambda}}| < \epsilon/2$ and (2.69) hold. We remark that \tilde{x}_∞ depends only on $\tilde{\lambda}, \tilde{P}, \epsilon$ since the number $\tilde{x}_{\infty,1}$ provided by Proposition 2.19 depends on θ_1 and θ_2 which depend only on $\tilde{\lambda}$ and ϵ .

By construction we have $\theta_1 < \theta_{\infty, \lambda} < \theta_2$ and (2.69) implies that the vector field of (2.25) points strictly outwards the cone between r_{θ_1} and r_{θ_2} for all $x \geq \tilde{x}_\infty$. Therefore, any nontrivial solution $z \in L^2(1, +\infty)$ of (2.25) approaches the origin at the angle $\theta_{\infty, \lambda}$ as x tends to infinity and a standard phase plane argument shows that $z(x)$ must remain inside the cone between r_{θ_1} and r_{θ_2} for all $x \geq \tilde{x}_\infty$. Hence (2.79) follows. \blacksquare

By means of the transformation $x = \phi_\beta(t)$, it is possible to prove an analogous result concerning the local behaviour of the angular coordinate when $x \rightarrow 0^+$; indeed, we have the following:

Proposition 2.22 *There exist $\epsilon_0 > 0$ such that for every $\epsilon \in (0, \epsilon_0)$, $\tilde{\lambda} \in \Lambda$ and $\tilde{P} \in \mathcal{P}_\mu$ there exist $\tilde{\delta}_0 > 0$ and $x_0 = x_0(\tilde{P}, \epsilon) > 0$ such that for every $\lambda \in \Lambda$ and $P \in \mathcal{P}_\mu$ with*

$$|\lambda - \tilde{\lambda}| < \tilde{\delta}_0, \quad \|P - \tilde{P}\|_{L^\infty(0,1)} < \tilde{\delta}_0 \quad (2.80)$$

and for every nontrivial solution $z \in L^2((0,1))$ of (2.1) we have

$$|\theta(x, \lambda) - \theta(0)| < \epsilon, \quad \forall x \in (0, x_0], \quad (2.81)$$

where $\theta(\cdot, \lambda)$ is any angular coordinate of z .

3 The linear eigenvalue problem

In this Section we are dealing with the study of the spectral theory for the linear operator formally defined by

$$\tau z = Jz' + P(x)z, \quad x > 0, \quad (3.1)$$

where $P \in \mathcal{P}_\mu$. Some information on the spectrum of τ follow directly from a standard spectral theory (see e.g. [18, 24]). Indeed, [24, Th. 6.8] ensures that τ is in the limit point case at infinity; moreover, from Remark 2.14 we deduce that τ is in the limit point case also at zero. Let us point out that this fact is a consequence of assumption (\mathcal{P}_3) on P^* .

Let us consider the operator A_0 defined by

$$\begin{aligned} D(A_0) &= \{z \in L^2(0, +\infty) : z \in AC(0, +\infty), \tau z \in L^2(0, +\infty)\}, \\ A_0 z &= \tau z, \quad \forall z \in D(A_0). \end{aligned} \quad (3.2)$$

From [24, Th. 5.8] we deduce that A_0 is the unique self-adjoint realization of τ ; moreover, arguing as in the proof of [18, Lemma 5.1], it is possible to see that $\sigma_{\text{ess}}(A_0) = (-\infty, \mu^-] \cup [\mu^+, +\infty)$.

As far as $D_0 := D(A_0)$ is concerned, we are able to prove the following result:

Proposition 3.1 *For every $z \in D_0$ we have*

$$z \in H^1(1, +\infty), \quad z \in L^\infty(0, +\infty).$$

PROOF. Assume that $z \in D_0$. Since $P \in L^\infty(1, +\infty)$ we deduce that $P(x)z \in L^2(1, +\infty)$; hence $Jz' = \tau z - P(x)z \in L^2(1, +\infty)$. This proves that $z \in H^1(1, +\infty) \subset L^\infty(1, +\infty)$.

The fact that $z \in L^\infty(0, 1)$ immediately follows from (2.59) and (2.60). ■

The aim of this Section is to study the problem of the existence of eigenvalues of A_0 in Λ ; first of all, let us observe that every eigenvalue of A_0 is simple, since τ is in the limit point case at infinity. Moreover, from Remark 2.14 we know that $\lambda \in \Lambda$ is an eigenvalue of A_0 if and only if there exists $c_\lambda \in \mathbb{R}$ such that

$$\zeta_{1,\lambda} = c_\lambda z_{1,\lambda}, \quad (3.3)$$

where $z_{1,\lambda}$ and $\zeta_{1,\lambda}$ are given in Proposition 2.4 and Proposition 2.11, respectively.

Remark 3.2 *According to Lemma 2.6 and Lemma 2.13, when $\lambda \in \Lambda$ is an eigenvalue of A_0 the associated eigenfunction z_λ satisfies $z_\lambda \in H_0^1(0, +\infty)$.*

In what follows we show that it is possible to write a condition equivalent to (3.3) by means of the angular function θ associated to solutions of (2.1) introduced in Subsection 2.2. To this aim, let us denote by $\vartheta(\cdot, \lambda)$ the angular coordinate of $\zeta_{1,\lambda}$, normalized in such a way that $\vartheta(0) \in (0, \pi)$, for every $\lambda \in \Lambda$.

From Proposition 2.16 we know that there exists

$$\lim_{x \rightarrow +\infty} \vartheta(x, \lambda) = \vartheta(+\infty, \lambda)$$

and that this limit corresponds to a function belonging to $H^1(1, +\infty)$ if and only if

$$\vartheta(+\infty, \lambda) = \pi - \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}} \pmod{\pi}. \quad (3.4)$$

Let us define $\nu : \Lambda \rightarrow \mathbb{R}$ by

$$\nu(\lambda) = \lim_{x \rightarrow +\infty} \vartheta(x, \lambda), \quad \forall \lambda \in \Lambda.$$

We then have the following characterization of the eigenvalues of A_0 :

Theorem 3.3 *A number $\lambda \in \Lambda$ is an eigenvalue of A_0 if and only if*

$$\nu(\lambda) = \pi - \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}} \pmod{\pi}. \quad (3.5)$$

In order to prove the existence of eigenvalues of A_0 it is then sufficient to study the behaviour of the function $\nu^* : \Lambda \rightarrow \mathbb{R}$ defined by

$$\nu^*(\lambda) = \nu(\lambda) + \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}}, \quad \forall \lambda \in \Lambda.$$

We will prove that ν^* is strictly increasing and continuous in Λ .

Proposition 3.4 *The function $\nu^* : \Lambda \rightarrow \mathbb{R}$ is strictly increasing in Λ .*

PROOF. Let us first observe that ν^* is the sum of ν and of the function ν_* defined by

$$\nu_*(\lambda) = \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}}, \quad \forall \lambda \in \Lambda;$$

since ν_* is strictly increasing in Λ , it is sufficient to prove that ν is increasing in Λ .

To this aim, let us recall (cf. [24, Cor. 16.2]) that for every fixed $x > 0$ the function

$$\begin{aligned} \varphi_x : \Lambda &\rightarrow \mathbb{R} \\ \lambda &\rightarrow \vartheta(x, \lambda) \end{aligned}$$

is increasing in Λ .

Now, let $\lambda, \lambda' \in \Lambda$ with $\lambda < \lambda'$; for every $x > 0$ we have

$$\vartheta(x, \lambda) \leq \vartheta(x, \lambda');$$

passing to the limit for $x \rightarrow +\infty$ we obtain

$$\lim_{x \rightarrow +\infty} \vartheta(x, \lambda) \leq \lim_{x \rightarrow +\infty} \vartheta(x, \lambda'),$$

i.e.

$$\nu(\lambda) \leq \nu(\lambda').$$

■

Proposition 3.5 *The function $\nu^* : \Lambda \rightarrow \mathbb{R}$ is continuous.*

PROOF. Let us observe again that it is sufficient to prove the continuity of ν . To this aim, let us fix $\tilde{\lambda} \in \Lambda$; let us consider $\epsilon > 0$ sufficiently small and apply Proposition 2.21 and Proposition 2.22 with $\tilde{P} = P$. Let $\delta_1 = \min(\tilde{\delta}, \tilde{\delta}_0)$ and let us denote by x_∞ and x_0 the numbers given in those Propositions.

Let us recall that a usual continuous dependence argument on the interval $[x_0, x_\infty]$, on which the equation (2.1) is not singular, proves that there exists $\delta_2 > 0$ such that if $|\lambda - \tilde{\lambda}| < \delta_2$ then

$$|(\vartheta(x_\infty, \lambda) - \vartheta(x_0, \lambda)) - (\vartheta(x_\infty, \tilde{\lambda}) - \vartheta(x_0, \tilde{\lambda}))| < \epsilon. \quad (3.6)$$

Consider now $\delta = \min(\delta_1, \delta_2)$ and assume that $|\lambda - \tilde{\lambda}| < \delta$; we can write

$$\begin{aligned} \nu(\lambda) - \nu(\tilde{\lambda}) &= \vartheta(+\infty, \lambda) - \vartheta(+\infty, \tilde{\lambda}) = \vartheta(+\infty, \lambda) - \vartheta(x_\infty, \lambda) + \\ &+ \vartheta(x_\infty, \lambda) - \vartheta(x_0, \lambda) + \vartheta(x_0, \tilde{\lambda}) - \vartheta(x_\infty, \tilde{\lambda}) + \vartheta(x_0, \lambda) - \vartheta(0, \lambda) + \\ &- \vartheta(x_0, \tilde{\lambda}) + \vartheta(0, \tilde{\lambda}) + \vartheta(x_\infty, \tilde{\lambda}) - \vartheta(+\infty, \tilde{\lambda}), \end{aligned} \quad (3.7)$$

taking into account that $\vartheta(0, \lambda) = \vartheta(0, \tilde{\lambda})$. From Proposition 2.21 and Proposition 2.22 we deduce that

$$\begin{aligned} |\vartheta(+\infty, \lambda) - \vartheta(x_\infty, \lambda)| &< \epsilon, \quad |\vartheta(x_\infty, \tilde{\lambda}) - \vartheta(+\infty, \tilde{\lambda})| < \epsilon \\ |\vartheta(x_0, \lambda) - \vartheta(0, \lambda)| &< \epsilon, \quad |\vartheta(x_0, \tilde{\lambda}) - \vartheta(0, \tilde{\lambda})| < \epsilon. \end{aligned} \quad (3.8)$$

From (3.6), (3.7) and (3.8) we obtain

$$|\nu(\lambda) - \nu(\tilde{\lambda})| < 5\epsilon$$

and this concludes the proof. ■

For every $k \in \mathbb{Z}$, let us denote by $\lambda_k \in \Lambda$ (if it exists) the number such that

$$\nu^*(\lambda_k) = k\pi,$$

i.e.

$$\vartheta(+\infty, \lambda_k) = k\pi + \pi - \arctan \sqrt{\frac{\lambda_k - \mu^-}{\mu^+ - \lambda_k}}. \quad (3.9)$$

The number λ_k is the ' k -th eigenvalue' of A_0 (if it exists) and we denote by $z_k \in D_0$ the corresponding eigenfunction; recalling (2.63), (2.65), from (3.9) and the fact that ν^* is strictly increasing we immediately deduce the following result:

Proposition 3.6 *For every $k \in \mathbb{Z}$ we have*

$$\begin{aligned} \text{rot}(z_k) &\in (k, k+1) & \text{if } \arctan \frac{w_{1,2}^*}{w_{1,1}^*} &\in (0, \pi/2) \\ \text{rot}(z_k) &\in (k-1/2, k+1/2) & \text{if } \arctan \frac{w_{1,2}^*}{w_{1,1}^*} &\in (\pi/2, \pi). \end{aligned} \quad (3.10)$$

Moreover, for every $k, l \in \mathbb{Z}$ with $k \neq l$ we also have

$$\text{rot}(z_k) \neq \text{rot}(z_l). \quad (3.11)$$

In what follows, we give some results on the accumulation of eigenvalues of A_0 at the boundary of Λ . We consider the (possible) accumulation at the end-point μ^+ ; conditions for accumulation at μ^- can be obtained in an analogous way.

From (3.5) and the definition of ν^* we infer that the existence of eigenvalues accumulating at μ^+ depends on the behaviour of ν^* in a left neighbourhood of μ^+ . This behaviour can be described by means of the limit

$$\lim_{\lambda \rightarrow (\mu^+)^-} \nu^*(\lambda), \quad (3.12)$$

whose existence is guaranteed from Proposition 3.4; more precisely, when the limit in (3.12) is infinite, then there exists $k_0 \in \mathbb{Z}$ such that for every $k \in \mathbb{Z}$, $k \geq k_0$, there exists $\lambda_k \in \Lambda$ for which (3.9) holds true and

$$\lim_{k \rightarrow +\infty} \lambda_k = \mu^+,$$

i.e. there is accumulation of eigenvalues at μ^+ . On the other hand, when the limit in (3.12) is finite, then there exists $M^+ \in \mathbb{R}$ such that

$$\nu^* \left(\frac{\mu^+ + \mu^-}{2} \right) < \nu^*(\lambda) < M^+, \quad \forall \lambda \in \left(\frac{\mu^+ + \mu^-}{2}, \mu^+ \right);$$

this implies that there is at most a finite number of eigenvalues of A_0 in $((\mu^+ + \mu^-)/2, \mu^+)$, i.e. there is not accumulation of eigenvalues at μ^+ .

Now, let us observe that the fact that the limit in (3.12) is finite or infinite depends on the analogous limit

$$\lim_{\lambda \rightarrow (\mu^+)^-} \nu(\lambda), \quad (3.13)$$

since the function ν_* is bounded in Λ . We are able to show that the finiteness of the limit in (3.13) depends on the behaviour of (2.1) when $\lambda = \mu^+$; to this aim, let us observe that a more careful analysis proves that Proposition 2.11 and Proposition 2.17 hold true also when $\lambda = \mu^+$. This implies that we are allowed to consider the solution $\zeta_{1,\lambda}$ of (2.1) with $\lambda = \mu^+$ satisfying (2.50) and the corresponding angular coordinate $\vartheta(\cdot, \mu^+)$, normalized in such a way that $\vartheta(0, \mu^+) \in (0, \pi)$.

Lemma 3.7 *Assume*

$$\lim_{x \rightarrow +\infty} \vartheta(x, \mu^+) = \theta^+ \in \mathbb{R}; \quad (3.14)$$

then we have

$$\lim_{\lambda \rightarrow (\mu^+)^-} \nu(\lambda) < +\infty. \quad (3.15)$$

PROOF. Let us observe that (3.14) implies that there exist $\Phi \in \mathbb{R}$ and $X > 0$ such that

$$\vartheta(x, \mu^+) < \Phi, \quad \forall x \geq X.$$

Moreover, from the monotonicity of $\vartheta(x, \cdot)$, for every $x \geq X$, we deduce that

$$\vartheta(x, \lambda) \leq \vartheta(x, \mu^+) < \Phi, \quad \forall \lambda < \mu^+.$$

Therefore, for every $\lambda \in \Lambda$ the function $\vartheta(\cdot, \lambda)$ is bounded from above by Φ in $[X, +\infty)$, hence we have

$$\nu(\lambda) = \lim_{x \rightarrow +\infty} \vartheta(x, \lambda) \leq \Phi, \quad \forall \lambda < \mu^+.$$

This is sufficient to conclude that (3.15) holds true. ■

Lemma 3.8 *Assume*

$$\lim_{x \rightarrow +\infty} \vartheta(x, \mu^+) = +\infty \quad (3.16)$$

and that there exists $X > 0$ such that

$$p_{11}(x) < \mu^-, \quad \forall x \geq X. \quad (3.17)$$

Then we have

$$\lim_{\lambda \rightarrow (\mu^+)^-} \nu(\lambda) = +\infty. \quad (3.18)$$

PROOF. Let us first observe that for every $\lambda \in \mathbb{R}$ the angular function $\vartheta(\cdot, \lambda)$ satisfies the differential equation

$$\theta' = (\lambda - p_{11}(x)) \cos^2 \theta - 2p_{12}(x) \cos \theta \sin \theta + (\lambda - p_{22}(x)) \sin^2 \theta. \quad (3.19)$$

From (3.19) and (3.17) we deduce that

$$\forall x \geq X, \lambda > \mu^- : \quad \vartheta(x, \lambda) = 0 \pmod{\pi} \quad \Rightarrow \quad \vartheta'(x, \lambda) > 0;$$

hence, if there exist $k \in \mathbb{Z}$ and $x_k \geq X$ such that

$$\vartheta(x_k, \lambda) > k\pi,$$

for some $\lambda > \mu^-$, then we can conclude that

$$\vartheta(x, \lambda) > k\pi, \quad \forall x \geq x_k.$$

Now, let us note that (3.16) implies that for every $M > 0$ there exists $x_M \geq X$ such that

$$\vartheta(x, \mu^+) > M + 2 + \pi, \quad \forall x \geq x_M$$

and let us fix $X^+ \geq x_M$; the continuity of $\vartheta(X^+, \cdot)$ ensures that there exists $\lambda_M < \mu^+$ such that

$$\vartheta(X^+, \lambda) > M + 1 + \pi, \quad \forall \lambda \in (\lambda_M, \mu^+).$$

According to the above remark, this implies that

$$\vartheta(x, \lambda) > M + 1, \quad \forall x \geq X^+, \lambda \in (\lambda_M, \mu^+)$$

and then

$$\nu(\lambda) = \lim_{x \rightarrow +\infty} \vartheta(x, \lambda) > M, \quad \forall \lambda \in (\lambda_M, \mu^+). \quad (3.20)$$

Therefore, for every $M > 0$ there exists $\lambda_M < \mu^+$ such that (3.20) holds, i.e.

$$\lim_{\lambda \rightarrow (\mu^+)^-} \nu(\lambda) = +\infty.$$

■

The question of the existence of eigenvalues can be dealt, arguing as in the proof of Proposition 3.18 in [4], as follows.

Proposition 3.9 *Assume that P has the form (2.8), where $\mu_a \in \mathbb{R}$, $k \in \mathbb{Z} \setminus \{0\}$ and $V \in C^1(0, +\infty)$ is a strictly increasing negative potential satisfying (2.9), with $\gamma_\infty < 0$ and $\alpha_\infty \in (0, 1]$, and (2.10). Then, the selfadjoint extension A_0 of the corresponding operator τ has a sequence of eigenvalues in $(-1, 1)$ accumulating at $\lambda = 1$.*

PROOF. We follow the same argument of [4, Prop. 3.15]. We first observe that the differential equation satisfied by $\vartheta(\cdot, 1)$ is

$$\vartheta'(x, 1) = 1 - \langle Q_{P(x)}[\cos \vartheta, \sin \vartheta], [\cos \vartheta, \sin \vartheta] \rangle,$$

where $Q_{P(x)}$ denotes the quadratic form associated to the matrix $P(x)$. By computing the eigenvalues of $P(x)$, we can prove that

$$\vartheta'(x, 1) \geq 1 - V(x) - \sqrt{1 + \left(\frac{k}{x} + \mu_a V'(x)\right)^2}, \quad \forall x \geq 1.$$

From assumption (2.9) we infer that

$$1 - V(x) - \sqrt{1 + \left(\frac{k}{x} + \mu_a V'(x)\right)^2} = -\frac{\gamma_\infty}{x^\alpha} + o\left(\frac{1}{x^\alpha}\right), \quad x \rightarrow +\infty;$$

this is sufficient to conclude that

$$\lim_{x \rightarrow +\infty} \vartheta(x, 1) = +\infty.$$

The result then follows from the application of Proposition 3.8. ■

A similar result (under more restrictive conditions on α) has been obtained by Schmid-Tretter in [18]; however, in [18] no information on the nodal properties of the eigenfunctions is provided.

4 The nonlinear eigenvalue problem

4.1 A bifurcation result

In this section we are interested in proving a global bifurcation result for a nonlinear equation of the form

$$\lambda z' + P(x)z = \lambda z + S(x, z)z, \quad \lambda \in \mathbb{R}, \quad x > 0, \quad z \in \mathbb{R}^2, \quad (4.1)$$

where $P \in \mathcal{P}_\mu$ and $S \in C((0, +\infty) \times \mathbb{R}^2, M_S^2)$. We denote by \mathcal{S} the set of continuous functions $S : (0, +\infty) \times \mathbb{R}^2 \rightarrow M_S^{2,2}$ satisfying the conditions

(\mathcal{S}_1) there exist $\alpha \in L^\infty(0, +\infty)$, $\eta_{ij} \in C(\mathbb{R}^2)$ such that $\eta_{ij}(0) = 0$, $i, j = 1, 2$, and

$$|S_{i,j}(x, z)| \leq \alpha(x)\eta_{ij}(z), \quad \forall x > 0, \quad z \in \mathbb{R}^2, \quad i, j = 1, 2; \quad (4.2)$$

(\mathcal{S}_2) for every compact $K \subset \mathbb{R}^2$ there exists $A_K > 0$ such that

$$\|S(x, z) - S(x, z')\| \leq A_K \|z - z'\|, \quad \forall x > 0, \quad z, z' \in K. \quad (4.3)$$

Let Σ denote the set of nontrivial solutions of (4.1) in $D_0 \times \Lambda$ and let $\Sigma' = \Sigma \cup \{(0, \lambda) \in D_0 \times \Lambda : \lambda \text{ is an eigenvalue of } A_0\}$, where D_0 and A_0 are as in Section 3. We denote by $\|\cdot\|_0$ the graph norm induced on D_0 by A_0 , defined as

$$\|z\|_0^2 = \|z\|_{L^2(0, +\infty)}^2 + \|\tau z\|_{L^2(0, +\infty)}^2, \quad \forall z \in D_0.$$

Let M denote the Nemitskii operator associated to S , given by

$$M(z)(x) = S(x, z(x))z(x), \quad \forall x > 0,$$

for every $z \in D_0$. We can show the validity of the following:

Proposition 4.1 *Assume that $S \in \mathcal{S}$ and that*

$$\lim_{x \rightarrow +\infty} \alpha(x) = 0, \tag{4.4}$$

where α is given in (4.2). Then $M : D_0 \longrightarrow L^2(0, +\infty)$ is a continuous compact map and satisfies

$$M(z) = o(\|z\|_0), \quad z \rightarrow 0. \tag{4.5}$$

The proof of Proposition 4.1 is based on the application of the following lemma:

Lemma 4.2 *Assume that $z_0, f_0 \in L^2(0, 1)$ satisfy*

$$\tau z_0 = f_0$$

and let $\{z_n\} \subset L^2(0, 1)$ be a sequence such that

$$\tau z_n = f_n,$$

for some $f_n \in L^2(0, 1)$. If $z_n \rightharpoonup z_0$ and $f_n \rightharpoonup f_0$ weakly in $L^2(0, 1)$, then $Mz_n \rightarrow Mz_0$ strongly in $L^2(0, 1)$.

PROOF. Let us apply Theorem 2.15 to the functions z_0 and z_n , for every $n \in \mathbf{N}$: we have

$$\begin{aligned} z_0(x) &= \nu_0(x) + w_0(x), \\ z_n(x) &= \nu_n(x) + w_n(x), \quad \forall x \in (0, 1), \end{aligned} \tag{4.6}$$

where

$$\nu_n(x) = c_n \zeta_1(x), \quad \nu_0(x) = c_0 \zeta_1(x)$$

$$w_n(x) = \int_0^1 G(x, \xi) f_n(\xi) d\xi, \quad w_0(x) = \int_0^1 G(x, \xi) f_0(\xi) d\xi, \quad \forall x \in (0, 1).$$

Since $G \in L^\infty((0, 1) \times (0, 1))$, we deduce that

$$w_n(x) \rightarrow w_0(x), \quad \forall x \in (0, 1) \tag{4.7}$$

by the weak convergence of f_n . Moreover, the estimate

$$\|w_n(x) - w_0(x)\| \leq \|G\|_{L^\infty((0, 1)^2)} \|f_n - f_0\|_{L^2(0, 1)}, \quad \forall x \in (0, 1), \tag{4.8}$$

holds true; the convergence $f_n \rightarrow f_0$ in $L^2(0,1)$ implies that the sequence $\{f_n\}$ is bounded in $L^2(0,1)$ and (4.7)-(4.8) ensure then that $w_n \rightarrow w_0$ in $L^2(0,1)$ by the dominated convergence theorem. This condition, together with the assumption $z_n \rightarrow z_0$ in $L^2(0,1)$, implies that $\nu_n \rightarrow \nu_0$ in $L^2(0,1)$. Hence, we obtain that $c_n \rightarrow c_0$, for $n \rightarrow +\infty$, and

$$\nu_n \rightarrow \nu_0 \quad \text{in } L^\infty(0,1) \text{ and in } L^2(0,1). \quad (4.9)$$

From (4.7)-(4.9) we have

$$z_n(x) \rightarrow z_0(x), \quad \forall x \in (0,1). \quad (4.10)$$

On the other hand, from (4.8) and the boundedness of $\{f_n\}$ in $L^2(0,1)$ we deduce also that $\{w_n\}$ is bounded in $L^\infty(0,1)$; as a consequence, using (4.9), we get that $\{z_n\}$ is bounded in $L^\infty(0,1)$. Using assumption (S_2) , from this fact we infer that there exists $C_1 > 0$ such that

$$\|S(x, z_n(x)) - S(x, z_0(x))\| \leq C_1 \|z_n(x) - z_0(x)\|, \quad \forall x \in (0,1); \quad (4.11)$$

equations (4.10)-(4.11) guarantee that

$$S(x, z_n(x)) \rightarrow S(x, z_0(x)) \quad \forall x \in (0,1). \quad (4.12)$$

Finally, from (4.11) and the boundedness of $\{z_n\}$ in $L^\infty(0,1)$ we also deduce that there exists $C_2 > 0$ such that

$$\|S(x, z_n(x))z_n(x) - S(x, z_0(x))z_0(x)\| \leq C_2, \quad \forall x \in (0,1), \quad \forall n \geq 1; \quad (4.13)$$

an application of the Lebesgue convergence Theorem gives

$$\int_0^1 \|S(x, z_n(x))z_n(x) - S(x, z_0(x))z_0(x)\|^2 dx \rightarrow 0, \quad n \rightarrow +\infty,$$

i.e. $Mz_n \rightarrow Mz$ in $L^2(0,1)$. ■

PROOF OF PROPOSITION 4.1. First of all, let us observe that it is sufficient to prove the result when $x \in (0,1)$. Indeed, the fact that $P \in L^\infty(1, +\infty)$ implies that the graph norm $\|\cdot\|_0$, when applied to functions defined on $[1, +\infty)$, is equivalent to the $H^1(1, +\infty)$ norm; hence, when $x \in [1, +\infty)$ we can apply [4, Prop. 4.3].

1. We first show that $Mz \in L^2(0, +\infty)$ when $z \in D_0$; from Proposition 3.1 we deduce that $z \in L^\infty(0,1)$. Therefore there exists $C_z > 0$ such that

$$|S(x, z(x))| \leq C_z, \quad \forall x \in (0,1).$$

As a consequence we obtain $Mz \in L^\infty(0,1) \subset L^2(0,1)$.

2. Let us fix $z_0 \in D_0$ and let $z_n \in D_0$ such that $z_n \rightarrow z_0$ when $n \rightarrow +\infty$; this implies that

$$z_n \rightarrow z_0 \text{ in } L^2(0,1), \quad \tau z_n \rightarrow \tau z_0 \text{ in } L^2(0,1) \quad (4.14)$$

We can then apply Lemma 4.2 and obtain that $Mz_n \rightarrow Mz_0$ in $L^2(0,1)$.

3. As far as the compactness of M is concerned, let $\{z_n\} \subset L^2(0,1)$ be such that

$$\|z_n\|_0 \leq K,$$

for some $K > 0$. This implies that, up to a subsequence, we have

$$z_n \rightharpoonup z_0 \text{ in } L^2(0, 1), \quad \tau z_n \rightharpoonup \tau z_0 \text{ in } L^2(0, 1). \quad (4.15)$$

Hence, according to Lemma 4.2, we conclude that $Mz_n \rightarrow Mz_0$ in $L^2(0, 1)$.

4. Finally, let us prove (4.5). We have

$$\|Mz\|_{L^2(0,1)}^2 = \int_0^1 \|Mz(x)\|^2 dz \leq \int_0^1 \|S(x, z(x))\|^2 \|z(x)\|^2 dx, \quad \forall z \in D_0. \quad (4.16)$$

Assume now that $z \rightarrow 0$ in D_0 ; this implies that $z \rightarrow 0$ and $\tau z \rightarrow 0$ in $L^2(0, 1)$; arguing as in the proof of Lemma 4.2, we deduce that $z \rightarrow 0$ in $L^\infty(0, 1)$. Therefore, assumption \mathcal{S}_2 implies that there exists $C > 0$ such that

$$\|S(x, z(x))\| \leq C\|z(x)\| \leq C\|z\|_{L^\infty(0,1)}, \quad \forall x \in (0, 1). \quad (4.17)$$

From (4.16) and (4.17) we deduce that

$$\|Mz\|_{L^2(0,1)} \leq C\|z\|_{L^\infty(0,1)}\|z\|_{L^2(0,1)} \leq C\|z\|_{L^\infty(0,1)}\|z\|_0,$$

which implies that $Mz = o(\|z\|_0)$ as $\|z\|_0 \rightarrow 0$. ■

Now, let us observe that, in view of the results on A_0 given in Section 3 and of Proposition 4.1, it is possible to write (4.1) as an abstract equation of the form

$$A_0 u + M(u) = \lambda u, \quad (u, \lambda) \in D_0 \times \mathbb{R}, \quad (4.18)$$

where $A_0 : D_0 \subset L^2(0, +\infty) \rightarrow L^2(0, +\infty)$ is an unbounded self-adjoint operator such that

$$\sigma_{ess}(A_0) = (-\infty, \mu^-] \cup [\mu^+, +\infty)$$

and $M : D_0 \times \mathbb{R} \rightarrow L^2(0, +\infty)$ is a continuous and compact map such that

$$M(u) = o(\|u\|), \quad u \rightarrow 0. \quad (4.19)$$

From an application of a global bifurcation result (see [22, Th. 1.2], [4, Th. 4.1]) to (4.18) we then obtain the following main result:

Theorem 4.3 *Assume that $P \in \mathcal{P}_\mu$, $S \in \mathcal{S}$ and that (4.4) holds true. Then, for every eigenvalue $\gamma \in \Lambda$ of A_0 there exists a continuum C_γ of nontrivial solutions of (4.1) in $D_0 \times \mathbb{R}$ bifurcating from $(0, \gamma)$ and such that one of the following conditions holds true:*

- (1) C_γ is unbounded in $D_0 \times \Lambda$;
- (2) $\sup\{\lambda : (u, \lambda) \in C_\gamma\} \geq \mu^+$ or $\inf\{\lambda : (u, \lambda) \in C_\gamma\} \leq \mu^-$;
- (3) C_γ contains $(0, \gamma') \in \Sigma'$, with $\gamma' \neq \gamma$.

Now, let us observe that a more precise description of the bifurcating branch, eventually leading to exclude condition (3), can be obtained when there exists a continuous functional $i : \Sigma' \rightarrow \mathbb{Z}$ (cf. [4, Th. 4.2]). In order to define such a functional, we first define the rotation number of solutions to (4.1) by means of a linearization procedure; to this aim for every solution (w, μ) of (4.1) we consider the linear equation

$$Jz' + P(x)z = \mu z + S(x, w(x))z, \quad (4.20)$$

which obviously reduces to

$$Jz' + P(x)z = \mu z \quad (4.21)$$

when $w = 0$. It is clear that w is a solution of (4.20); let us denote by P_w the matrix defined by

$$P_w(x) = P(x) - S(x, w(x)), \quad \forall x > 0.$$

We can prove the following result:

Lemma 4.4 *For every $(w, \mu) \in \Sigma$ we have $P_w \in \mathcal{P}_\mu$.*

PROOF. Let us first observe that $w \in D_0$ implies that $w \in H^1(1, +\infty)$ and $w \in L^\infty(0, +\infty)$ (cf. Proposition 3.1). In particular we have

$$\lim_{x \rightarrow +\infty} w(x) = 0 \quad (4.22)$$

and there exists a compact set $K_w \subset \mathbb{R}^2$ such that

$$w(x) \in K_w, \quad \forall x > 0.$$

Using (2.4), assumption (\mathcal{S}_1) and the fact that $w \in L^\infty(0, 1)$, we obtain that

$$\lim_{x \rightarrow 0^+} x^\beta P_w(x) = \lim_{x \rightarrow 0^+} (x^\beta P(x) - x^\beta S(x, w(x))) = P^*;$$

therefore P_w satisfies (2.4). Moreover, we have

$$R_{0,w}(x) = x^\beta P_w(x) - P^* = R_0(x) - x^\beta S(x, w(x)), \quad \forall x > 0. \quad (4.23)$$

Using again (\mathcal{S}_1) , we plainly deduce that there exists $\eta \in C(\mathbb{R}^2, \mathbb{R}^+)$ such that

$$\int_0^1 \frac{1}{x^\beta} \|x^\beta S(x, w(x))\|^{q_0} dx = \int_0^1 x^{\beta(q_0-1)} \alpha(x) \eta(w(x)) dx < +\infty, \quad (4.24)$$

since $q_0 \geq 1$ and $\alpha, w \in L^\infty(0, 1)$. From (2.5), (4.23) and (4.24) we can conclude that $R_{0,w}$ satisfies (2.5).

Now, we pass to the proof of the validity of (\mathcal{P}_1) . Using (2.2), assumption (\mathcal{S}_1) and (4.22), we infer that

$$\lim_{x \rightarrow +\infty} P_w(x) = \lim_{x \rightarrow +\infty} (P(x) - S(x, w(x))) = P_\infty;$$

hence P_w satisfies (2.2).

Moreover, we have

$$R_{\infty,w}(x) = P_w(x) - P_\infty = R_\infty(x) - S(x, w(x)), \quad \forall x > 0. \quad (4.25)$$

From assumption (S_2) , with $K = K_w$ and $z' = 0$, we obtain

$$\int_1^{+\infty} \|S(x, w(x))\|^2 dx \leq A_{K_w}^2 \int_1^{+\infty} \|w(x)\|^2 dx < +\infty. \quad (4.26)$$

When $q_\infty \geq 2$ this allows to conclude that $R_{\infty,w}$ satisfies (2.3), since

$$\begin{aligned} \int_1^{+\infty} \|S(x, w(x))\|^{q_\infty} dx &= \int_1^{+\infty} \|S(x, w(x))\|^{q_\infty-2} \|S(x, w(x))\|^2 dx \leq \\ &\leq C_w \int_1^{+\infty} \|S(x, w(x))\|^2 dx < +\infty. \end{aligned} \quad (4.27)$$

Finally, also when $q_\infty < 2$ it is possible to show that $R_{\infty,w}$ satisfies (2.3) with the same q_∞ of R_∞ ; indeed, at this point we can say that $w \in H^1(1, +\infty)$ is a nontrivial solution of the linear equation

$$Jz' + P_w(x)z = \mu z,$$

where $P_w \in \mathcal{P}_\mu$ and $\mu \in \Lambda$. Therefore Proposition 2.4 applies and we deduce that w satisfies the first condition in (2.26). As a consequence, $w \in L^{q_\infty}(1, +\infty)$ and we are able to repeat (4.26) with the exponent q_∞ instead of 2. \blacksquare

As a consequence of Lemma 4.4, the results of Section 2 apply to (4.20); in particular, when $w \neq 0$ we can consider the number $\text{rot}(w)$ defined in (2.65).

Definition 4.5 Assume that $P \in \mathcal{P}_\mu$ and $S \in \mathcal{S}$ and let (w, μ) be a solution of (4.1).

If $(w, \mu) \neq (0, \mu)$, then the rotation number $j(w, \mu)$ of (w, μ) is defined by

$$j(w, \mu) = \text{rot}(w). \quad (4.28)$$

If $(w, \mu) = (0, \mu)$ and the linear problem (4.21) has a nontrivial solution z_μ belonging to $H^1(0, +\infty)$, then the rotation number $j(w, \mu)$ of (w, μ) is defined by

$$j(w, \mu) = \text{rot}(z_\mu). \quad (4.29)$$

By means of Definition 4.5 we have defined $j : \Sigma' \rightarrow \mathbb{R}$; this functional will be used in order to construct a continuous discrete functional whose values are preserved in the bifurcating branches C_γ of solutions of (4.1). It is important now to observe that every branch C_γ satisfies

$$C_\gamma \subset H_0^1(0, +\infty) \times \mathbb{R};$$

indeed, this is a consequence that $(z, \lambda) \in C_\gamma$ is a solution of the linear equation

$$Jz' + P_z(x)z = \lambda z$$

such that $z \in D_0$. According to Remark 3.2 this implies that $z \in H_0^1(0, +\infty)$.

Hence, it is sufficient to study the continuity properties of j with respect to the $H_0^1(0, +\infty)$ -norm, denoted by $||| \cdot |||$.

Proposition 4.6 *The function $j : \Sigma' \rightarrow \mathbb{R}$ is continuous.*

PROOF. We prove the continuity of j at every point $(w, \mu) \in \Sigma \cap H_0^1(0, +\infty)$. In a very similar way it is possible to show that j is also continuous at every point $(0, \lambda)$, with λ eigenvalue of A_0 .

Let us fix $(w, \mu) \in \Sigma \cap H_0^1(0, +\infty)$ and let $\epsilon > 0$ small enough; consider then the numbers δ , δ_0 , x_∞ and x_0 given in Proposition 2.21 and Proposition 2.22 (with $\tilde{\lambda} = \mu$ and $\tilde{P} = P_w$) and let $\delta_1 = \min(\delta, \delta_0)$.

Using assumption (S_1) and the continuous embedding $H_0^1(0, +\infty) \subset L^\infty(0, +\infty)$, it is possible to show that there exist $\delta_2 > 0$ such that

$$||P_z - P_w||_{L^\infty(0, +\infty)} < \delta_1$$

if $|||z - w||| < \delta_2$.

Hence, from Proposition 2.21 and Proposition 2.22 we deduce that for every $(z, \lambda) \in \Sigma \cap H_0^1(0, +\infty)$ with $|\lambda - \mu| < \delta_2$ and $|||z - w||| < \delta_2$ we have

$$\begin{aligned} |\theta_z(x, \lambda) - \theta_z(+\infty, \lambda)| &< \epsilon, \quad \forall x \geq x_\infty \\ |\theta_z(x, \lambda) - \theta_z(0)| &< \epsilon, \quad \forall x \in (0, x_0]. \end{aligned} \tag{4.30}$$

Now, let us observe that we have

$$j(z, \lambda) - j(w, \mu) = \frac{\theta_z(+\infty, \lambda) - \theta_z(0)}{\pi} - \frac{\theta_w(+\infty, \mu) - \theta_w(0)}{\pi} = \frac{\theta_z(+\infty, \lambda) - \theta_w(+\infty, \mu)}{\pi},$$

since $\theta_z(0) = \theta_w(0)$. Therefore, the result follows from the same argument used in the proof of Proposition 3.5. \blacksquare

Before defining the functional i a remark is in order; we recall that if $(w, \lambda) \in \Sigma \cap H_0^1(0, +\infty)$ then

$$\theta_w(0) = \arctan \frac{w_{1,2}^*}{w_{1,1}^*} \in (0, \pi) \mod \pi$$

and

$$\theta_w(+\infty, \lambda) = \pi - \arctan \sqrt{\frac{\lambda - \mu^-}{\mu^+ - \lambda}} \in \left(\frac{\pi}{2}, \pi\right) \mod \pi.$$

As a consequence, when

$$\arctan \frac{w_{1,2}^*}{w_{1,1}^*} \in \left(0, \frac{\pi}{2}\right) \tag{4.31}$$

we have

$$j(w, \lambda) \notin \mathbb{Z}, \quad \forall (w, \lambda) \in \Sigma \cap H_0^1(0, +\infty).$$

On the other hand, if

$$\arctan \frac{w_{1,2}^*}{w_{1,1}^*} \in \left(\frac{\pi}{2}, \pi \right) \quad (4.32)$$

we have

$$j(w, \lambda) + \frac{1}{2} \notin \mathbb{Z}, \quad \forall (w, \lambda) \in \Sigma \cap H_0^1(0, +\infty).$$

This suggests to define $i : \Sigma' \rightarrow \mathbb{Z}$ as

$$i(w, \lambda) = [j(w, \lambda)], \quad \forall (w, \lambda) \in \Sigma', \quad (4.33)$$

if (4.31) holds true, and

$$i(w, \lambda) = \left[j(w, \lambda) + \frac{1}{2} \right], \quad \forall (w, \lambda) \in \Sigma', \quad (4.34)$$

if (4.32) holds true (recall also Proposition 3.6). Let us observe that Proposition 3.6 also implies that

$$i(z_\gamma, 0) \neq i(z_{\gamma'}, 0), \quad (4.35)$$

for every $\gamma \neq \gamma' \in \Lambda$ eigenvalues of A_0 (with associated eigenfunctions z_γ and $z_{\gamma'}$, respectively).

From Proposition 4.6 and the definition of i we obtain the following result:

Proposition 4.7 *The function $i : \Sigma' \rightarrow \mathbb{R}$ is continuous.*

As a consequence, using Proposition 4.7 and (4.35), from Theorem 4.3 we deduce the final result:

Theorem 4.8 *Assume that $P \in \mathcal{P}_\mu$, $S \in \mathcal{S}$ and that (4.4) hold true. Then, for every eigenvalue $\gamma \in \Lambda$ of A_0 there exists a continuum C_γ of nontrivial solutions of (4.1) in $D_0 \times \mathbb{R}$ bifurcating from $(0, \gamma)$ and such that one of the conditions (1)-(2) of Theorem 4.3 holds true and*

$$i(w, \lambda) = i(z_\gamma, 0), \quad \forall (w, \lambda) \in C_\gamma, \quad (4.36)$$

where z_γ is the eigenfunction of A_0 associated to γ .

4.2 Application to the Dirac equation

Let us consider the partial differential equation

$$i \sum_{j=1}^3 \alpha_j \frac{\partial \psi}{\partial x_j} - \beta \psi - V(\|x\|) \psi + i a \sum_{j=1}^3 \alpha_j \frac{\partial V(\|x\|)}{\partial x_j} \psi = \lambda \psi + \gamma(\|x\|) F(\langle \beta \psi, \psi \rangle) \beta \psi, \quad x \in \mathbb{R}^3, a \in \mathbb{R}, \quad (4.37)$$

where $\psi : \mathbb{R}^3 \rightarrow \mathbb{C}^4$, $V \in C((0, +\infty), \mathbb{R})$ satisfies (2.9)-(2.10)-(2.11)-(2.12), $\gamma \in C((0, +\infty), \mathbb{R})$ fulfills

$$\lim_{r \rightarrow 0^+} r^2 \gamma(r) \in \mathbb{R}, \quad r^2 \gamma(r) = o(1), \quad r \rightarrow +\infty, \quad (4.38)$$

$F \in C(\mathbb{R}, \mathbb{R})$, $\langle \cdot, \cdot \rangle$ denotes the scalar product in \mathbb{C}^4 and α_j ($j = 1, 2, 3$) and β are the 4×4 matrices given by

$$\alpha_j = \begin{pmatrix} 0 & \sigma_j \\ \sigma_j & 0 \end{pmatrix}, \quad \beta = \begin{pmatrix} \sigma_0 & 0 \\ 0 & -\sigma_0 \end{pmatrix},$$

where

$$\sigma_0 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \quad \sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

We remark that nonlinearities like the one in (4.37) give rise to the so-called generalized Soler models (see [14]). In fact, Soler [21] formulated a model of extended fermions by introducing a self interaction term which corresponds to the choice $F(s) = s$ in (4.37) (see [17] for a survey on interaction terms which are interesting from a physical point of view).

We denote by H_0 the (free) Dirac operator defined by

$$H_0\psi = i \sum_{j=1}^3 \alpha_j \frac{\partial \psi}{\partial x_j} - \beta \psi, \quad \forall \psi \in H_0^1(\mathbb{R}^3) \subset L^2(\mathbb{R}^3). \quad (4.39)$$

In [23] a decomposition of $H_0 - V + ia \alpha \cdot \nabla V$ has been performed, using polar coordinates in \mathbb{R}^3 and the unitary isomorphism

$$\begin{aligned} \varphi : L^2(\mathbb{R}^3) &\rightarrow L^2((0, +\infty), dr; L^2(S^2)) \\ \psi &\mapsto \tilde{\psi}, \end{aligned} \quad (4.40)$$

where $\tilde{\psi}$ is defined by

$$\tilde{\psi}(r, \theta, \phi) = r\psi(x(r, \theta, \phi)), \quad \forall r > 0, (\theta, \phi) \in S^2. \quad (4.41)$$

In order to describe such a decomposition, for every $l = 0, 1, 2, \dots$ and $m = -l, -l+1, \dots, l$ let us denote by Y_l^m the usual spherical harmonic; moreover, for every $j = 1/2, 3/2, 5/2, \dots$, let $m_j = -j, -j+1, \dots, j$ and $k_j = -(j+1/2), j+1/2$ and define

$$\Psi_{j-1/2}^{m_j} = \frac{1}{\sqrt{2j}} \begin{pmatrix} \sqrt{j+m_j} Y_{j-1/2}^{m_j-1/2} \\ \sqrt{j-m_j} Y_{j-1/2}^{m_j+1/2} \end{pmatrix}, \quad \Psi_{j+1/2}^{m_j} = \frac{1}{\sqrt{2j+2}} \begin{pmatrix} \sqrt{j+1-m_j} Y_{j+1/2}^{m_j-1/2} \\ -\sqrt{j+1+m_j} Y_{j+1/2}^{m_j+1/2} \end{pmatrix} \quad (4.42)$$

and

$$\Phi_{m_j, \mp(j+1/2)}^+ = \begin{pmatrix} i \Psi_{j \mp 1/2}^{m_j} \\ 0 \end{pmatrix}, \quad \Phi_{m_j, \mp(j+1/2)}^- = \begin{pmatrix} 0 \\ \Psi_{j \pm 1/2}^{m_j} \end{pmatrix}. \quad (4.43)$$

We also set

$$\mathcal{H}_{m_j, k_j} = \text{span} (\Phi_{m_j, k_j}^+, \Phi_{m_j, k_j}^-), \quad \forall j = 1/2, 3/2, \dots \quad (4.44)$$

Then, we have the following result:

Theorem 4.9 ([23, Th. 4.14]) *For every $j = 1/2, 3/2, \dots$ the subspace $C_0^\infty(0, +\infty) \otimes \mathcal{H}_{m_j, k_j} \subset L^2((0, +\infty), dr; L^2(S^2)^4)$ is invariant under the action of $H_0 - V + ia \alpha \cdot \nabla V$. Moreover, with respect to the basis $\{\Phi_{m_j, k_j}^+, \Phi_{m_j, k_j}^-\}$ of \mathcal{H}_{m_j, k_j} the restriction of $H_0 - V + ia \alpha \cdot \nabla V$ to \mathcal{H}_{m_j, k_j} can be represented by the operator h_{m_j, k_j} given by*

$$h_{m_j, k_j} = \begin{pmatrix} -1 & -V & \frac{d}{dr} - \frac{k_j}{r} + aV' \\ -\frac{d}{dr} - \frac{k_j}{r} + aV' & 1 - V \end{pmatrix}. \quad (4.45)$$

Moreover, the Dirac operator $H_0 - V + ia \alpha \cdot \nabla V$ on $C_0^\infty(\mathbb{R}^3)^4$ is unitarily equivalent to the direct sum of the partial waves operators h_{m_j, k_j} , i.e.

$$H_0 - V + ia \alpha \cdot \nabla V \approx \bigoplus_{j=1/2, 3/2, \dots}^{+\infty} \bigoplus_{m_j=-j}^j \bigoplus_{k_j=\pm(j+1/2)} h_{m_j, k_j}.$$

Remark 4.10 *The partial wave subspaces can be considered as a suitable generalization of radial functions adapted to the structure of the nonlinear problem. More precisely, the vectors Φ_{m_j, k_j}^\pm which are a basis for the partial wave subspace \mathcal{H}_{m_j, k_j} are the eigenfunctions of the spin orbit operator (cf. ([23])). We also observe that these subspaces are implicitly used in [1], [2], where (having in mind the Soler model) the system of ODEs is obtained from the PDE by making the ansatz that solutions should be a linear combination of functions of the form $\Phi_{1/2, 1}^+ = (\frac{i}{2\sqrt{\pi}} \cos \theta, \frac{i}{2\sqrt{\pi}} e^{i\phi} \sin \theta, 0, 0)$, $\Phi_{1/2, 1}^- = (0, 0, \frac{1}{2\sqrt{\pi}}, 0)$. On the same lines but in the context of the Schrödinger equation, we refer to [24, Example 1.5].*

Let us observe that the operators $\tau_{k_j} = h_{m_j, k_j}$, $j = 1/2, 3/2, \dots$, are of the form (3.1) with $P = P_{V, k_j, a}$ as in (2.8). Therefore, we can apply the theory developed in Sections 2 and 3; in particular we can consider the selfadjoint realization A_0 of τ_{k_j} , $j = 1/2, 3/2, \dots$, defined in (3.2). We denote by A_{k_j} this operator, by D_{k_j} its domain and we define

$$E_{k_j} = \{u^+ \Phi_{m_j, k_j}^+ + u^- \Phi_{m_j, k_j}^- : u = (u^+, u^-) \in D_{k_j}\}.$$

From Theorem 4.9 and the definition of D_{k_j} we immediately deduce that the image of E_{k_j} via the operator $H_0 - V + ia \alpha \cdot \nabla V$ is contained in $L^2((0, \infty)) \otimes \mathcal{H}_{k_j, m_j}$, for every $j = 1/2, 3/2, \dots$

Now, let us observe that Theorem 4.9 states that the subspaces

$$C_0^\infty(0, +\infty) \otimes \mathcal{H}_{m_j, k_j} \subset L^2((0, +\infty), dr; L^2(S^2)^4), \quad j = 1/2, 3/2, \dots,$$

are preserved by the linear operator in (4.37). It is important to note that in the particular case of $j = 1/2$ the subspaces $C_0^\infty(0, +\infty) \otimes \mathcal{H}_{m_{1/2}, k_{1/2}} \subset L^2((0, +\infty), dr; L^2(S^2)^4)$ are invariant also for the nonlinear term $F(\langle \beta \psi, \psi \rangle) \beta \psi$ in (4.37) (cf. [3, Lemma 5.5]), when F is regular.

Indeed, let $u \in L^2(\mathbb{R}^3)^4$ such that $\varphi(u) \in C_0^\infty(0, +\infty) \otimes \mathcal{H}_{m_{1/2}, k_{1/2}}$, where φ is defined in (4.40)-(4.41). A simple computation, based on the expressions of the functions $\Phi_{m_{1/2}, k_{1/2}}^\pm$, shows that, if we have

$$(\varphi(u))(r, \theta, \phi) = u^+(r) \Phi_{m_{1/2}, k_{1/2}}^+(\theta, \phi) + u^-(r) \Phi_{m_{1/2}, k_{1/2}}^-(\theta, \phi),$$

then

$$\langle \beta u(x), u(x) \rangle = \frac{1}{4\pi^2 r^2} [(u^+(r))^2 - (u^-(r))^2] \quad (4.46)$$

and

$$\begin{aligned} & F \left(\left\langle \beta \frac{(\varphi(u))(r, \theta, \phi)}{r}, \frac{(\varphi(u))(r, \theta, \phi)}{r} \right\rangle \right) \beta [(\varphi(u))(r, \theta, \phi)] \\ &= F \left(\frac{1}{4\pi^2 r^2} (u^+(r))^2 - (u^-(r))^2 \right) \left(u^+(r) \Phi_{m_{1/2}, k_{1/2}}^+(\theta, \phi) - u^-(r) \Phi_{m_{1/2}, k_{1/2}}^-(\theta, \phi) \right), \end{aligned} \quad (4.47)$$

proving that

$$\tilde{\psi} \in C_0^\infty(0, +\infty) \otimes \mathcal{H}_{m_{1/2}, k_{1/2}} \Rightarrow F \left(\left\langle \beta \frac{\tilde{\psi}}{r}, \frac{\tilde{\psi}}{r} \right\rangle \right) \beta \tilde{\psi} \in C_0^\infty(0, +\infty) \otimes \mathcal{H}_{m_{1/2}, k_{1/2}}. \quad (4.48)$$

Then, with an argument similar to the one developed in the proof of Proposition 4.1, we deduce that

$$\tilde{\psi} \in E_{k_{1/2}} \Rightarrow \gamma(r) F \left(\left\langle \beta \frac{\tilde{\psi}}{r}, \frac{\tilde{\psi}}{r} \right\rangle \right) \beta \tilde{\psi} \in L^2(0, +\infty) \otimes \mathcal{H}_{m_{1/2}, k_{1/2}}. \quad (4.49)$$

This fact is important to obtain a relation between solutions of (4.37) and solutions of a nonlinear ordinary differential equation of the form (4.1). Indeed, for every function $u \in L^2(\mathbb{R}^3)^4$ with $\varphi(u) \in E_{k_{1/2}}$, let $z = (u^+, u^-) \in D_{k_{1/2}}$ such that

$$\varphi(u) = u^+ \Phi_{m_{1/2}, k_{1/2}}^+ + u^- \Phi_{m_{1/2}, k_{1/2}}^-. \quad (4.50)$$

Then, (4.49) implies that $u \in L^2(\mathbb{R}^3)^4$ with $\varphi(u) \in E_{k_{1/2}}$ is a nontrivial solution of (4.37) if and only if $z = (u^+, u^-) \in D_{k_{1/2}}$ is a nontrivial solution of

$$\tau_{k_{1/2}} z = \lambda z + \gamma(r) F \left(\frac{(u^+)^2 - (u^-)^2}{4\pi r^2} \right) \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} z, \quad r > 0. \quad (4.51)$$

Let us denote $E = \varphi^{-1}(E_{k_{1/2}})$; in view of the above arguments and choosing

$$S(r, z) = \gamma(r) F \left(\frac{(u^+)^2 - (u^-)^2}{4\pi r^2} \right) \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix},$$

from Theorem 4.8 we plainly obtain the following result:

Theorem 4.11 *Let us suppose that $V \in C(0, +\infty)$ and $\gamma \in C(0, +\infty)$ satisfy (2.9)-(2.10)-(2.11) and (4.38) and let $F : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ be a locally Lipschitz continuous function such that $|F(s)| \leq C|s|$ for all $s \in \mathbb{R}^2$ and some constant $C > 0$. Then, for every eigenvalue $\mu \in (-1, 1)$ of $A_{k_{1/2}}$ there exists a continuum C_μ of nontrivial solutions of (4.37) in $E \times \mathbb{R}$ such that one of the conditions*

- (1) C_μ is unbounded in $E \times (-1, 1)$
- (2) $\sup\{\lambda : (u, \lambda) \in C_\mu\} \geq 1$ or $\inf\{\lambda : (u, \lambda) \in C_\mu\} \leq -1$

holds true and

$$\tilde{i}(w, \lambda) = i(z_\mu, 0), \quad \forall (w, \lambda) \in C_\mu, \quad (4.52)$$

where

$$\tilde{i}(w, \lambda) = i((w^+, w^-), \lambda)$$

and z_μ is the eigenfunction of $A_{k_{1/2}}$ associated to μ .

References

- [1] BALABANE M., CAZENAVE T., DOUADY A. and MERLE F., Existence of excited states for a non-linear Dirac field, *Commun. Math. Phys.*, **119** (1988), 153–176.
- [2] BALABANE M., CAZENAVE T. and VAZQUEZ L., Existence of standing waves for Dirac fields with singular nonlinearities, *Commun. Math. Phys.*, **133** (1990), 53–74.
- [3] CACCIAFESTA F., Global small solutions to the critical radial Dirac equation with potential, *Non-linear Anal.*, **74** (2011), 6060–6073.
- [4] CAPIETTO A. and DAMBROSIO W., Planar Dirac-type systems: the eigenvalue problem and a global bifurcation result, *J. London Math. Soc.*, **81** (2010), 477–498.
- [5] CAPIETTO A., DAMBROSIO W. and PAPINI D., A global bifurcation result for a second order singular equation, *Rend. Istit. Mat. Univ. Trieste, Special Issue in honour of Prof. F. Zanolin*, **44** (2012), 173–185.
- [6] COPPEL W.A., *Dichotomies in stability theory*, Lectures Notes in Mathematics **629**, 1978.
- [7] DANCER N., Boundary-value problems for ordinary differential equations on infinite intervals, *Proc. London Math. Soc.* **30** (1975), 76–94.
- [8] DANCER N., Boundary-value problems for ordinary differential equations of infinite intervals. II, *Quart. J. Math. Oxford Ser.* **28** (1977), 101–115.
- [9] DING J., XU, J. and ZHANG, F., Solutions of super-linear Dirac equations with general potentials, *Differ. Equ. Dyn. Syst.*, **17** (2009), 235–256.
- [10] DING Y. and RUF, B., Solutions of a nonlinear Dirac equation with external fields, *Arch. Ration. Mech. Anal.*, **190** (2008), 57–82.
- [11] DONG Y. and XIE, J., Admissible solutions for Dirac equations with singular and non-monotone nonlinearity, *Proc. Edinb. Math. Soc.*, **54** (2011), 363–371.
- [12] DUNFORD N. and SCHWARTZ J., *Linear Operators - Part II: spectral theory*, Interscience Publishers, 1963.
- [13] EASTHAM M.S.P., *The asymptotic solution of linear differential systems*, London Math. Society Monographs New Series, 1989.
- [14] ESTEBAN M.J., An overview on linear and nonlinear Dirac equations, *Discrete Contin. Dyn. Syst.*, **8** (2002), 381–397.

- [15] KALF H. and SCHMIDT K.M., Spectral stability of the Coulomb-Dirac Hamiltonian with anomalous magnetic moment, *J. Differential Equations* **205** (2004), 408–423.
- [16] RABIER P.J. and STUART C., Global bifurcation for quasilinear elliptic equations on \mathbb{R}^N , *Math. Z.*, **237** (2001), 85–124.
- [17] RAÑADA A.F., *On nonlinear classical Dirac fields and quantum physics*, in Old and New Questions in Physics, Cosmology, Philosophy, and Theoretical Biology, Plenum, New York, 1983, pp. 363–376.
- [18] SCHMID H. and TRETTER C., Singular Dirac systems and Sturm-Liouville problems nonlinear in the spectral parameter, *J. Differential Equations*, **181** (2002), 511–542.
- [19] SCHMID H. and TRETTER C., Eigenvalue accumulation for Dirac operators with spherically symmetric potential, *J. Phys. A*, **37** (2004), 8657–8674.
- [20] SECCHI S. and STUART C., Global bifurcation of homoclinic solutions of Hamiltonian systems, *Discrete Contin. Dyn. Syst.*, **9** (2003), 1493–1518.
- [21] SOLER M., Classical, Stable, Nonlinear Spinor Field with Positive Rest Energy, *Phys. Rev. D*, **1** (1970), 2766–2769.
- [22] STUART C., Global properties of components of solutions of non-linear second order differential equations on the half-line, *Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4)* **2** (1975), 265–286.
- [23] THALLER B., *The Dirac Equation*, Text and Monographs in Physics, 1992.
- [24] WEIDMANN J., *Spectral theory of ordinary differential equations*, Lectures Notes in Mathematics **1258**, 1987.